

3 Species Fermion Gases

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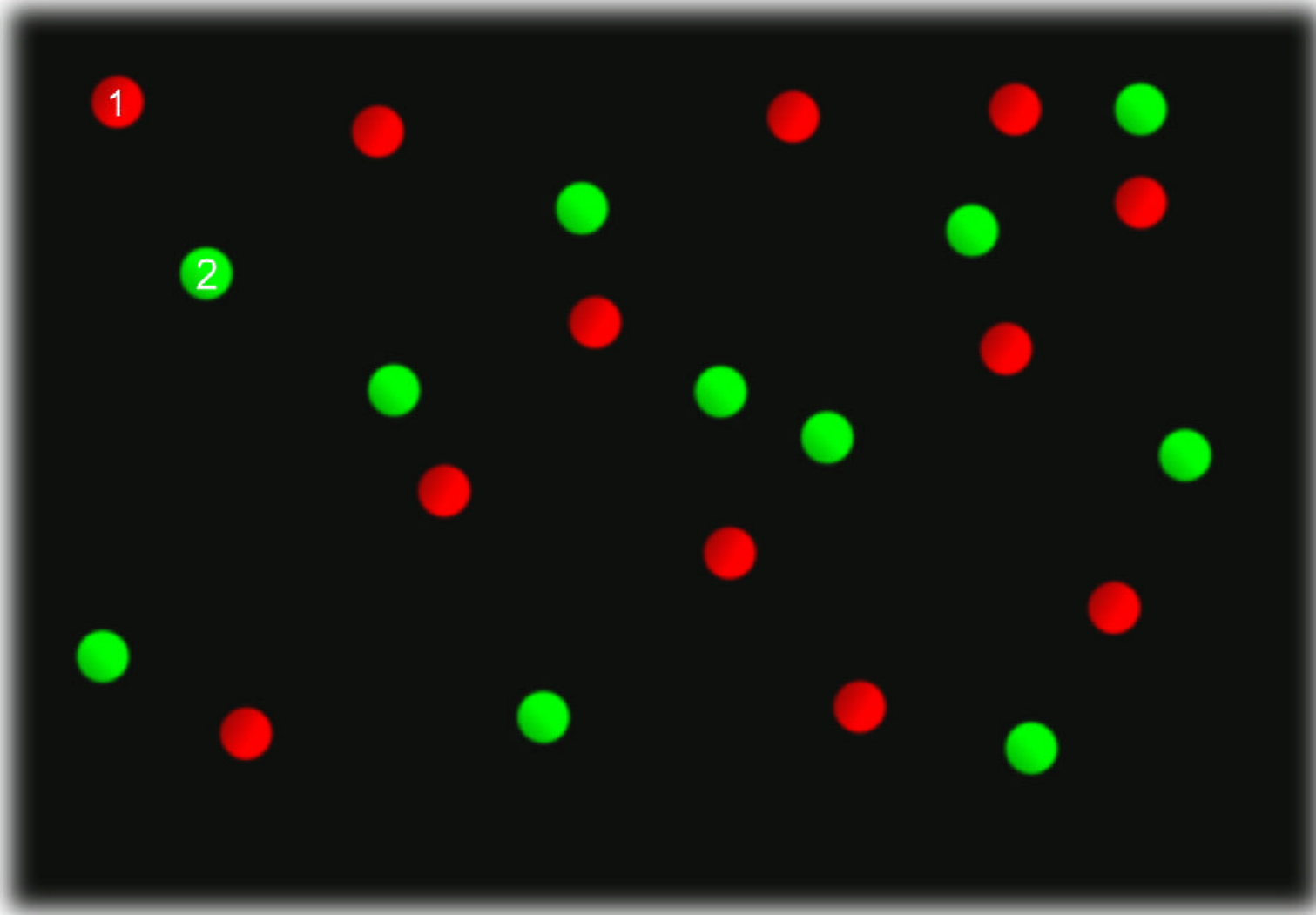
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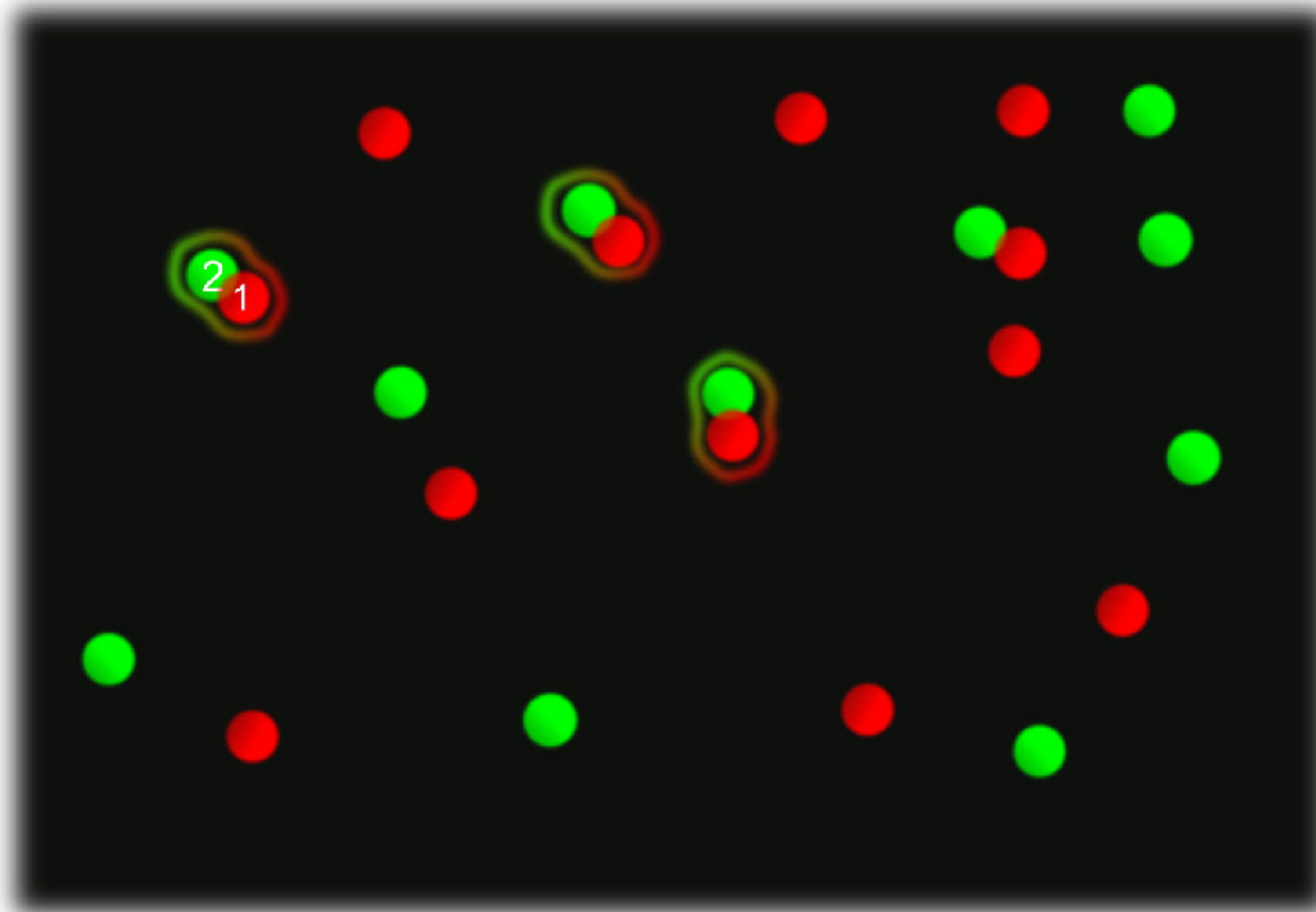
Goal:

Understanding of three fermion species systems

First: Consider a free 2 component system (no interactions)



Now add attractive interaction: Binding to bosonic states might happen

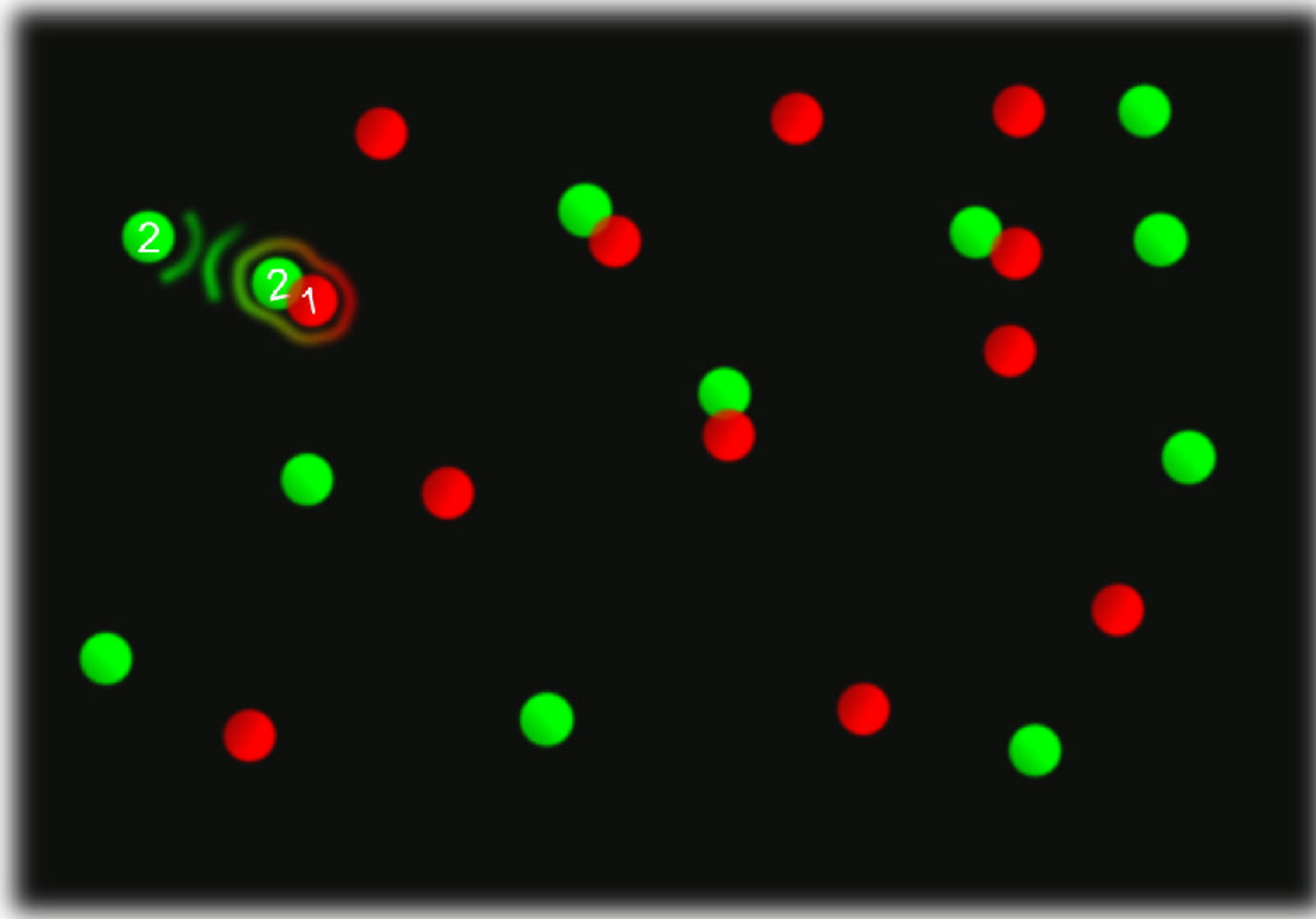


What's about a three fermion bound state?

Pauli blocking prohibits formation of such a state.

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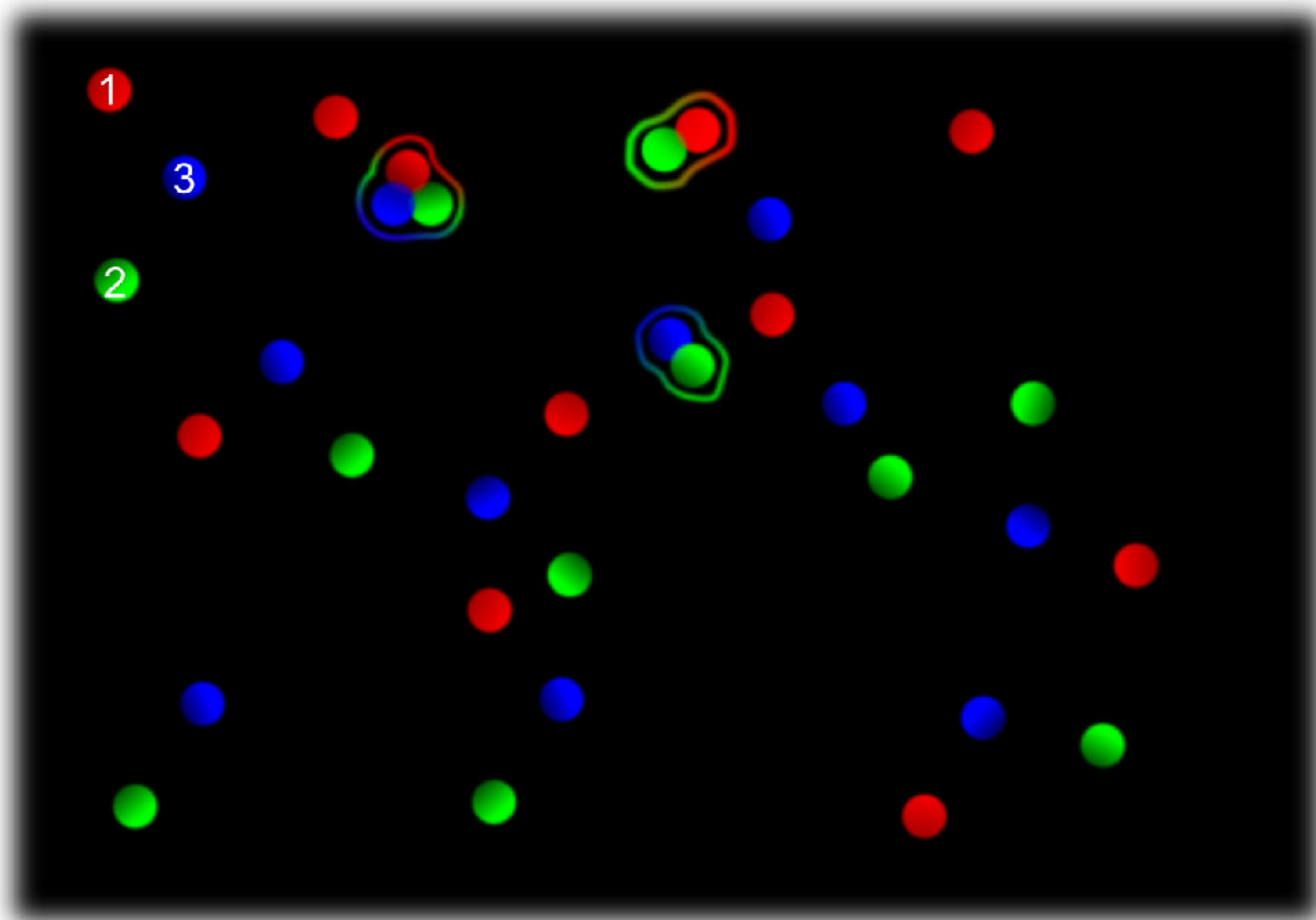
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Introduction

Let's add a third species of fermions

In principle the formation of a three fermion bound state (the *trion*) is possible.



The Question:
Do these 3 fermion bound states exist?

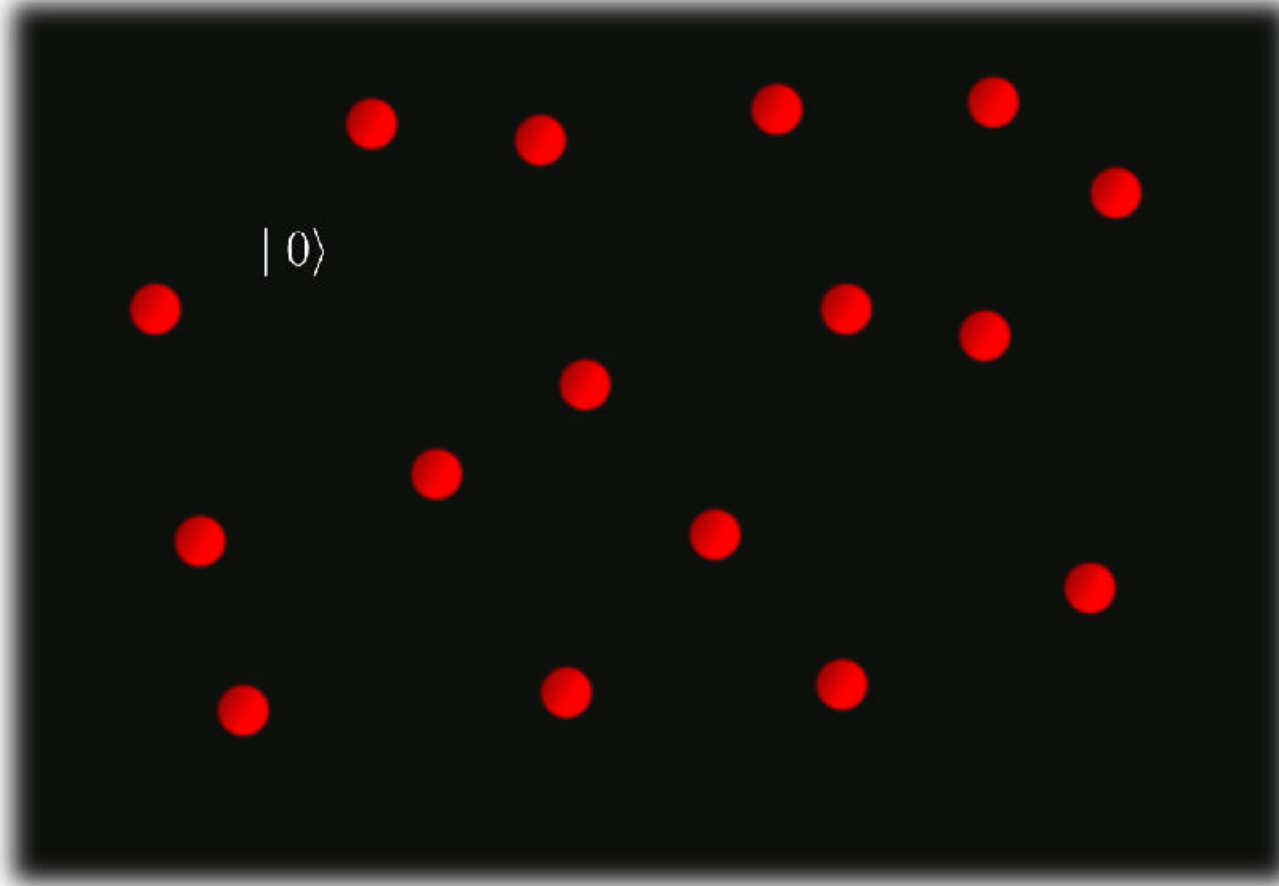
Remark: For bosonic systems they are shown to exist: *Efimov states*

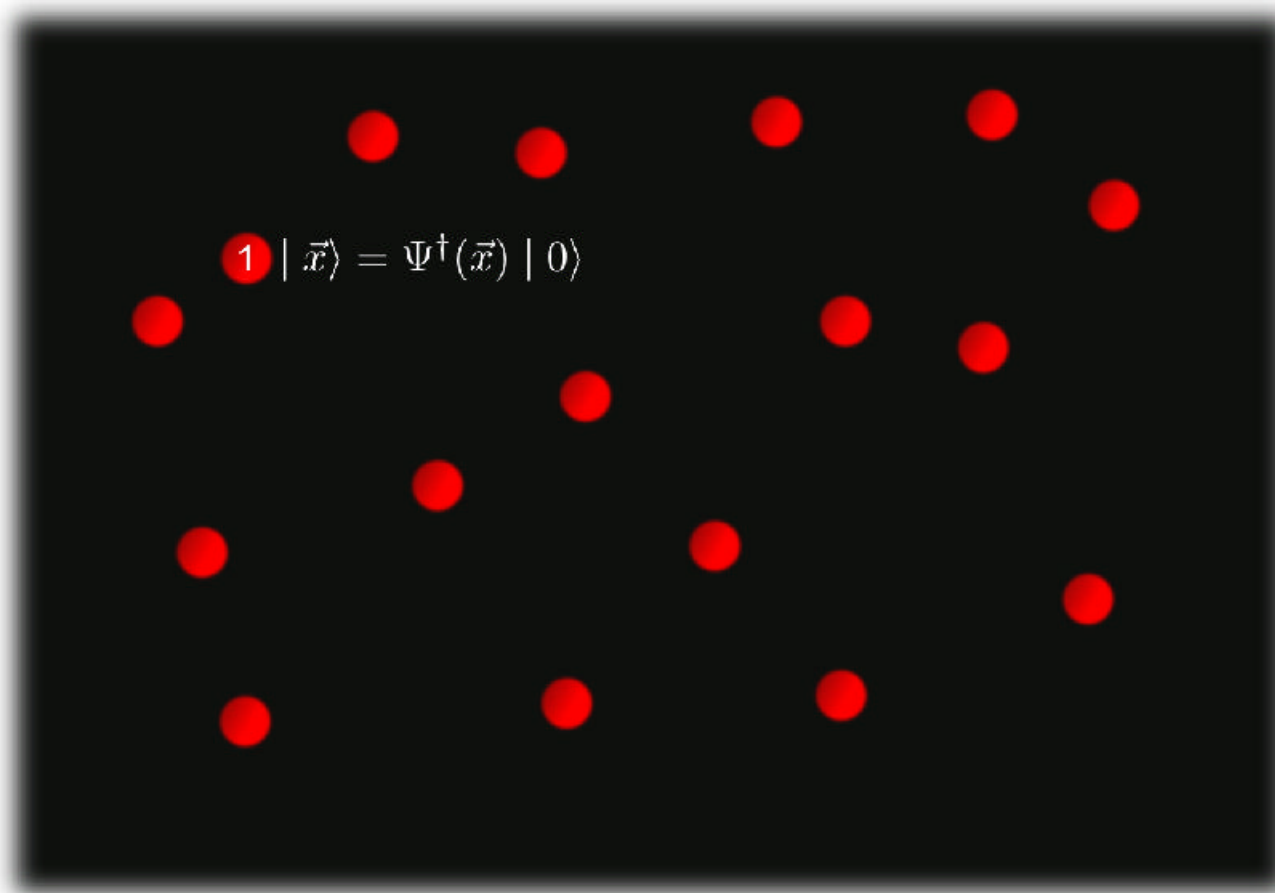
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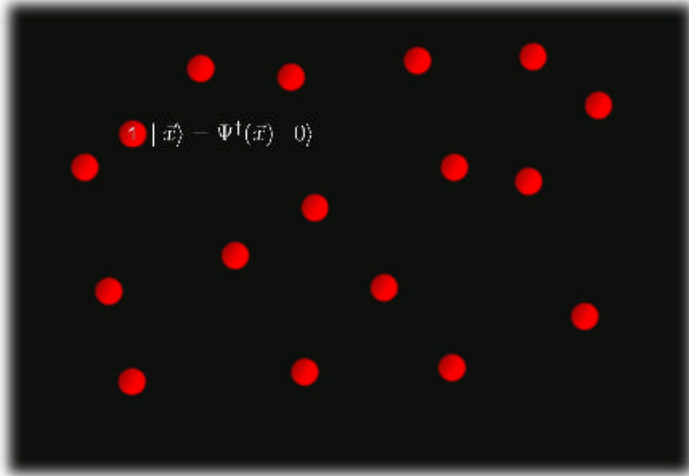
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.1 Getting familiar with the Hamiltonian

.1.1 The one component gas







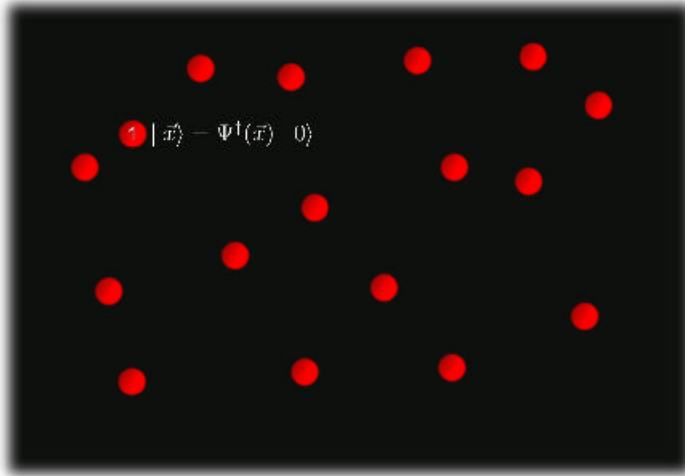
$\hat{\Psi}^\dagger(\vec{x})$ applied to the vacuum creates a new particle at position x .

$\hat{\Psi}(\vec{x})$ annihilates a particle at position x .

$|0\rangle$ denotes the vacuum.

$|\vec{x}\rangle$ represents a particle at position x .

How does the Hamiltonian for such a system look like?



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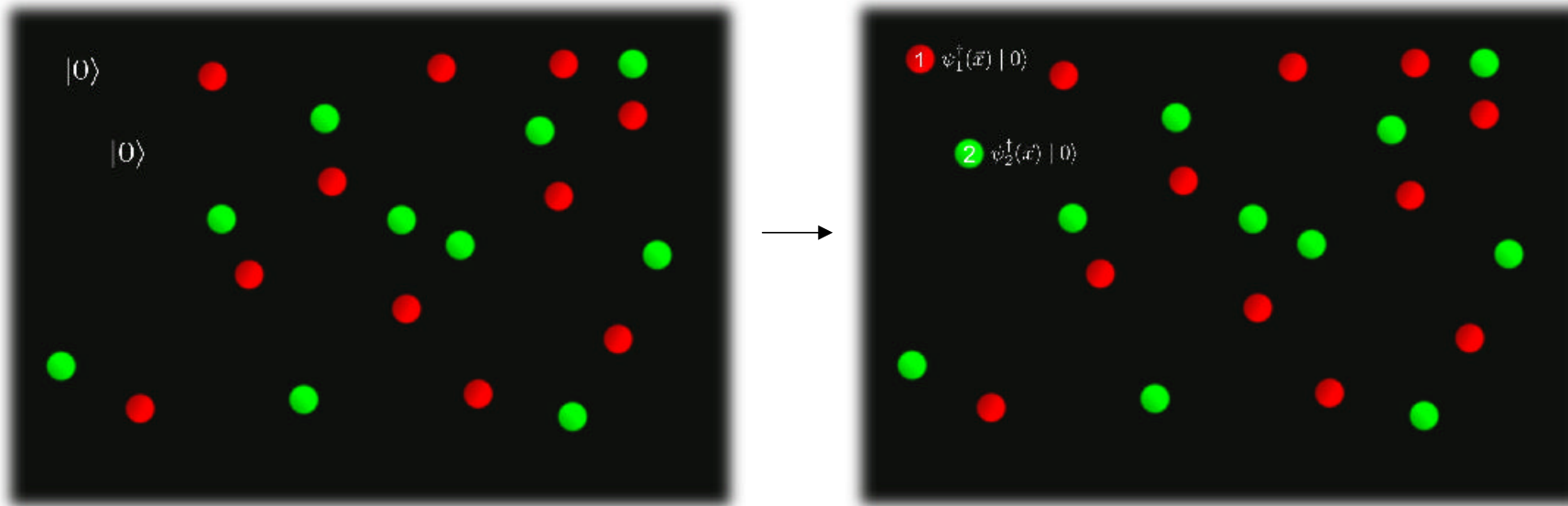
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How does the Hamiltonian for such a system look like?

$$\hat{H} = \int d^3x \hat{\psi}^\dagger(\vec{x}) \left(-\frac{\hbar^2}{2m} \Delta + V(\vec{x}) \right) \hat{\psi}(\vec{x})$$

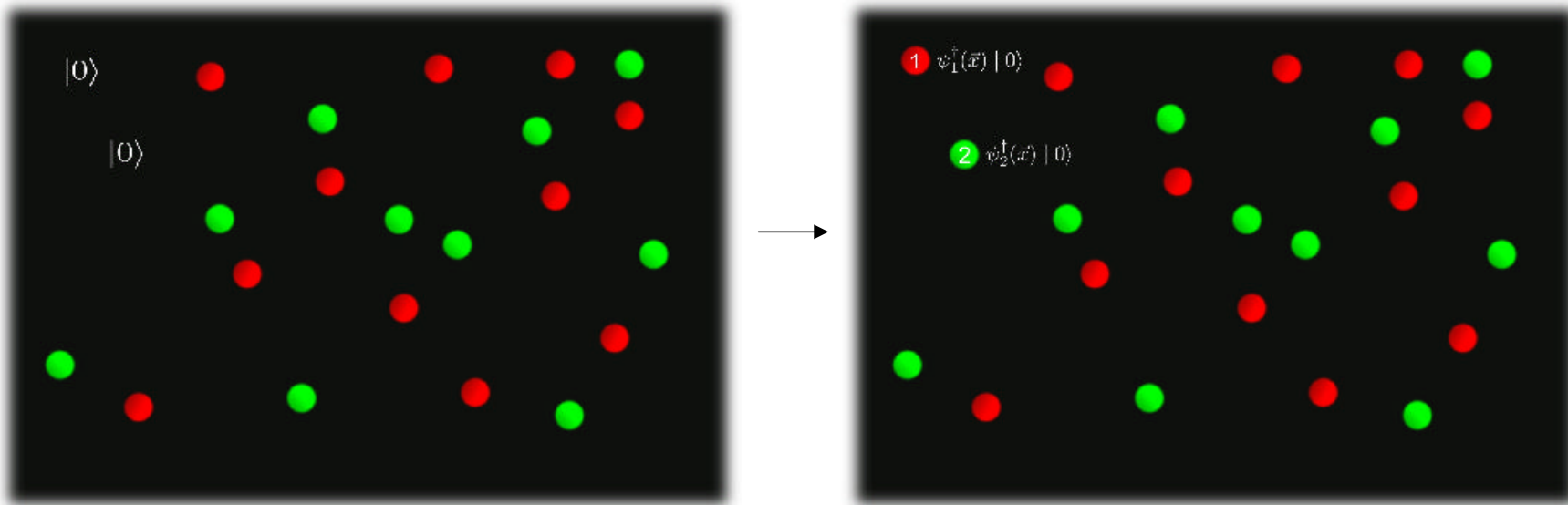
2.1.2 Two Component Gas



We need one more field in the Hamiltonian. Let's label the fields with $\sigma = 1, 2$:

$$\hat{H} = \sum_{\sigma=1,2} \int d^3x \hat{\psi}_{\sigma}^{\dagger}(\vec{x}) \left(-\frac{\hbar^2}{2m_{\sigma}} \Delta \right) \hat{\psi}_{\sigma}(\vec{x})$$

But now there is also the possibility of interactions between the different fermion species. How do we incorporate this?



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1.3 Interactions in 2-component systems

Interactions between two particles are described by a scattering process.

There are different possible scattering processes between two fermions with an attractive potential:

- (a) Elastic scattering
- (b) Inelastic scattering (leads to molecule formation)

In diagrammatic language process (a) is described by

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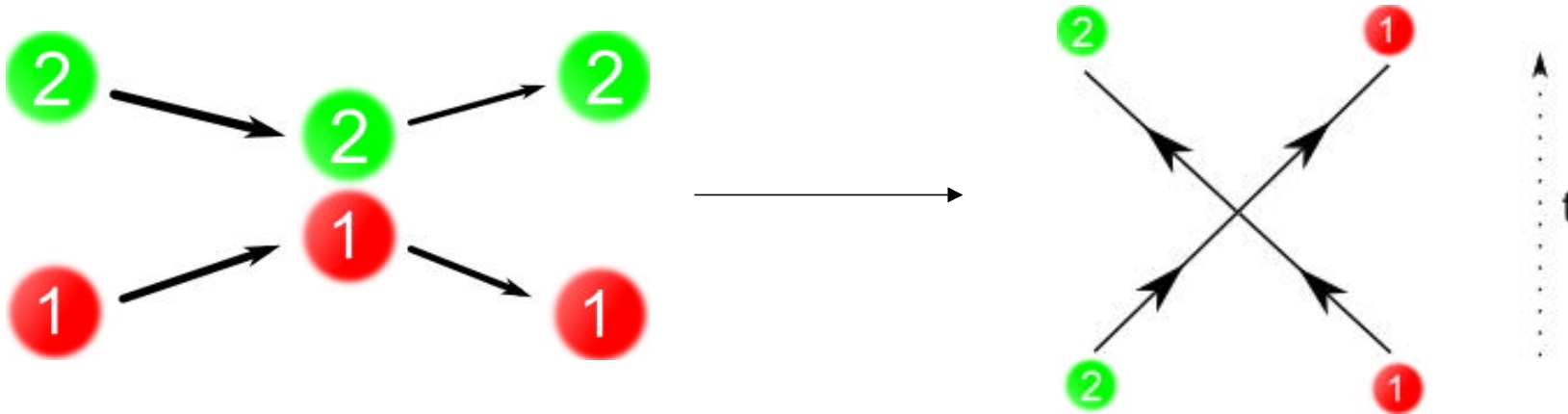
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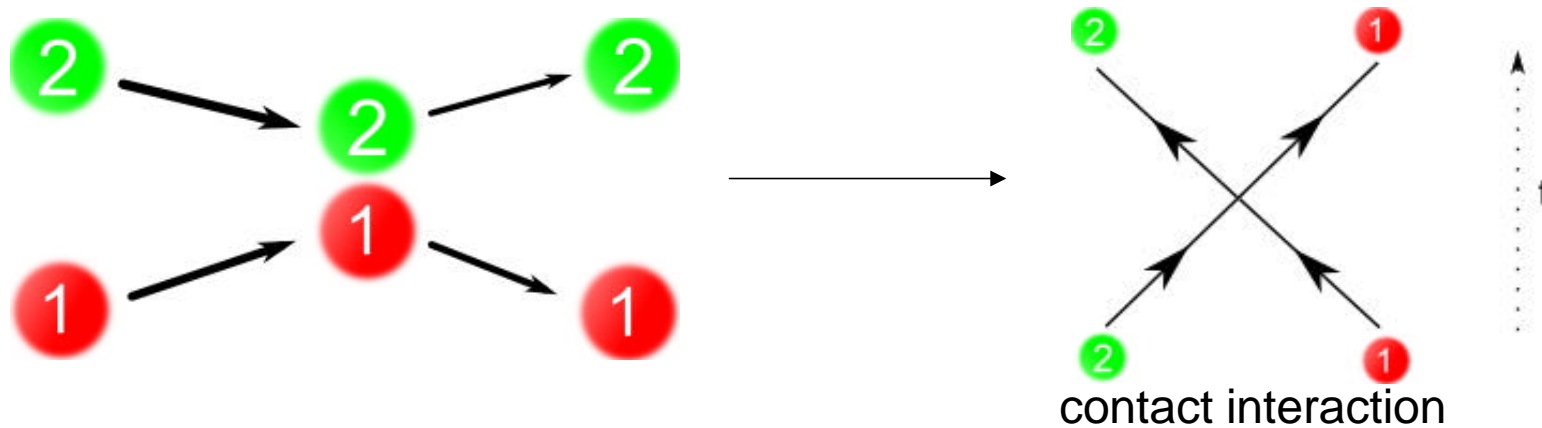
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In the Hamiltonian this interaction corresponds to a term of the form

$$\lambda_{12} \hat{\psi}_1^\dagger(\vec{x}) \hat{\psi}_2^\dagger(\vec{x}) \hat{\psi}_2(\vec{x}) \hat{\psi}_1(\vec{x})$$

$\hat{\psi}_i(\vec{x})$ describes an incoming particle.

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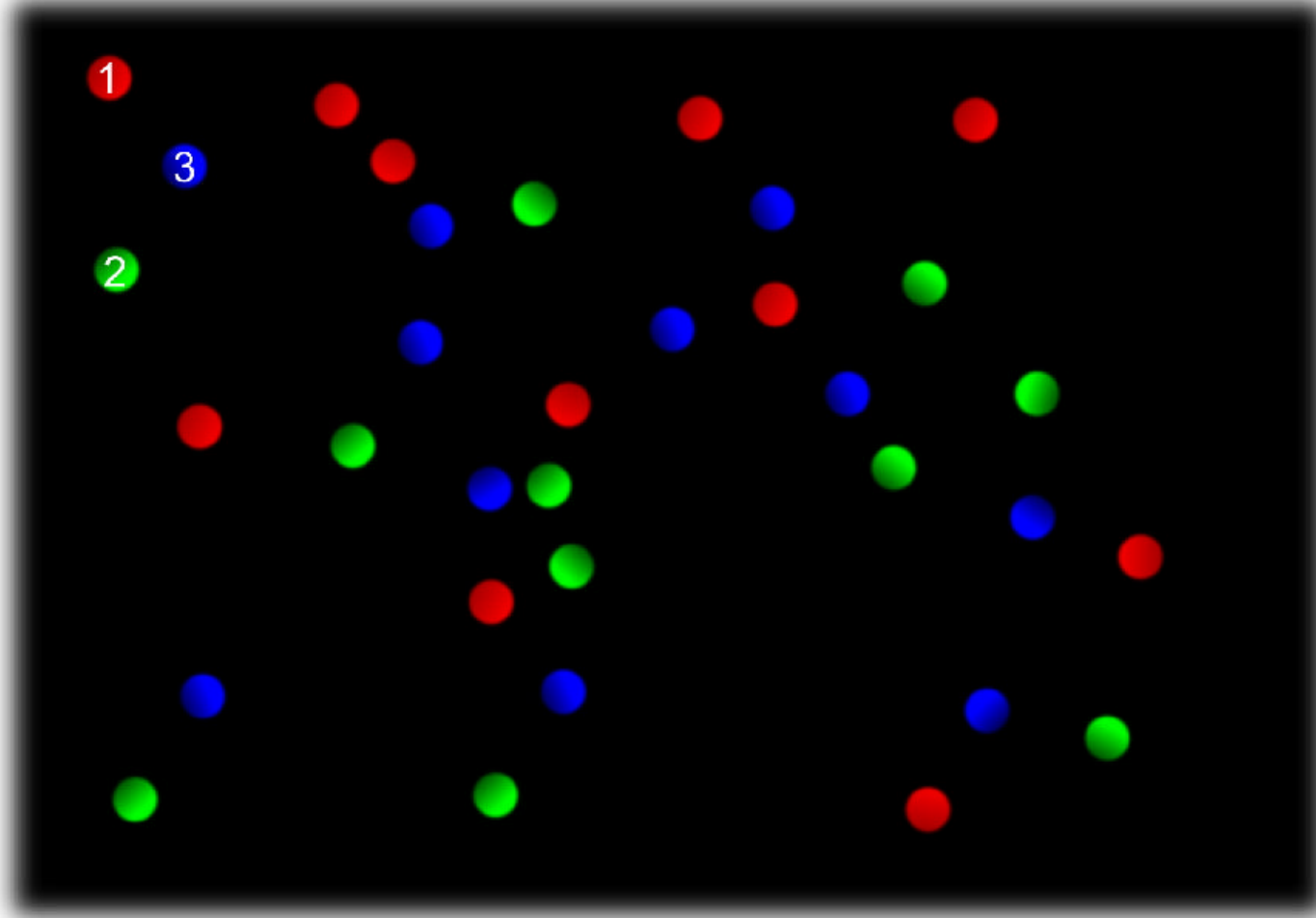
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$$\begin{aligned} \hat{H} = & \sum_{\sigma=1,2} \int d^3x \hat{\psi}_\sigma^\dagger(\vec{x}) \left(-\frac{\hbar^2}{2m_\sigma} \Delta \right) \hat{\psi}_\sigma(\vec{x}) \\ \rightarrow & + \int d^3x \lambda_{12} \hat{\psi}_1^\dagger(\vec{x}) \hat{\psi}_2^\dagger(\vec{x}) \hat{\psi}_2(\vec{x}) \hat{\psi}_1(\vec{x}) \end{aligned}$$

.2 The Three Fermion system

Now we add the third fermion species



What do we need?

- One additional field operator: $\hat{\psi}_3(\vec{x}), \hat{\psi}_3^\dagger(\vec{x})$
- Possible scattering between all combinations of fermions, thus we have 3 interaction terms:

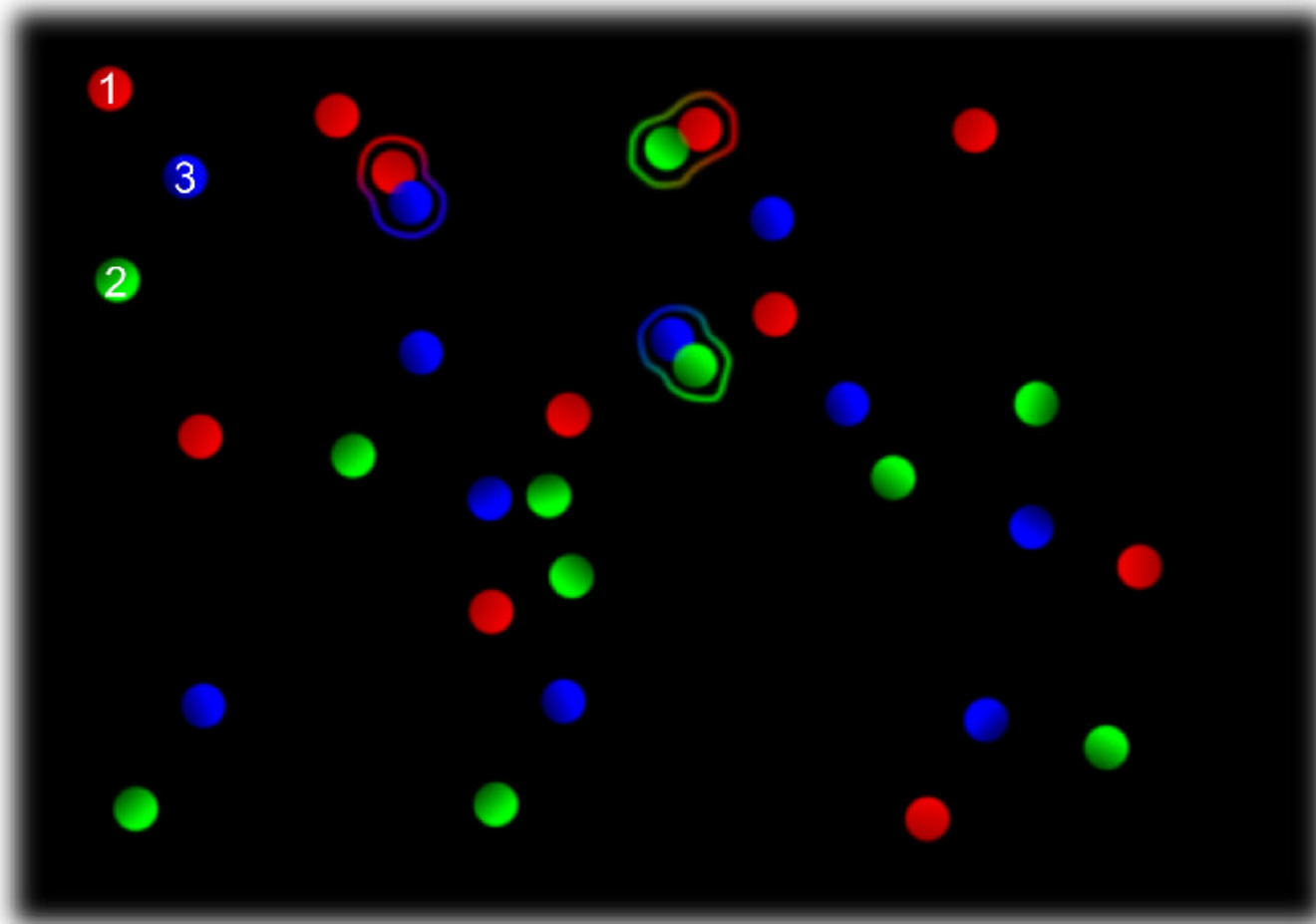
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 & + \int d^3x \lambda_{12} \hat{\psi}_1^\dagger(\vec{x}) \hat{\psi}_2^\dagger(\vec{x}) \hat{\psi}_2(\vec{x}) \hat{\psi}_1(\vec{x}) \longrightarrow \text{Diagram 1} \\
 & + \int d^3x \lambda_{13} \hat{\psi}_1^\dagger(\vec{x}) \hat{\psi}_3^\dagger(\vec{x}) \hat{\psi}_3(\vec{x}) \hat{\psi}_1(\vec{x}) \longrightarrow \text{Diagram 2} \\
 & + \int d^3x \lambda_{23} \hat{\psi}_2^\dagger(\vec{x}) \hat{\psi}_3^\dagger(\vec{x}) \hat{\psi}_3(\vec{x}) \hat{\psi}_2(\vec{x}) \longrightarrow \text{Diagram 3}
 \end{aligned}$$

.3 A Little on BEC-BCS Theory, Statistical Field Theory, Symmetry Breaking and all that

Where are the molecules?



Recall

- (a) Elastic scattering
- (b) Inelastic scattering (leads to molecule formation)

This time we are interested in process (b), describing molecule formation.

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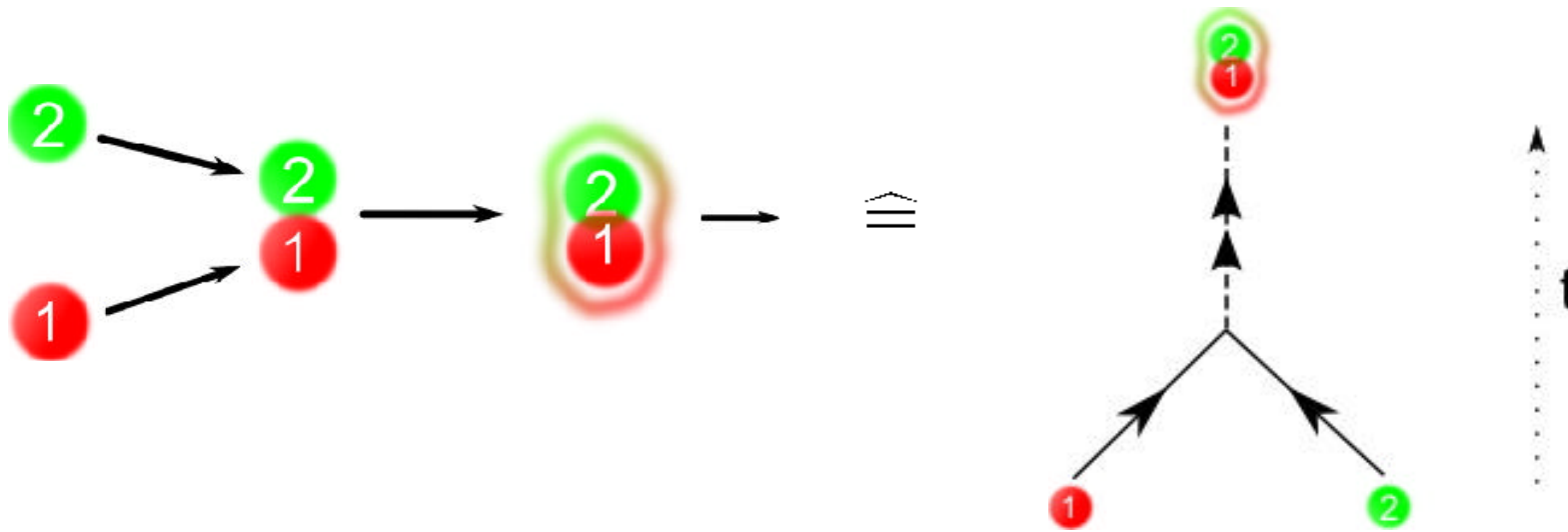
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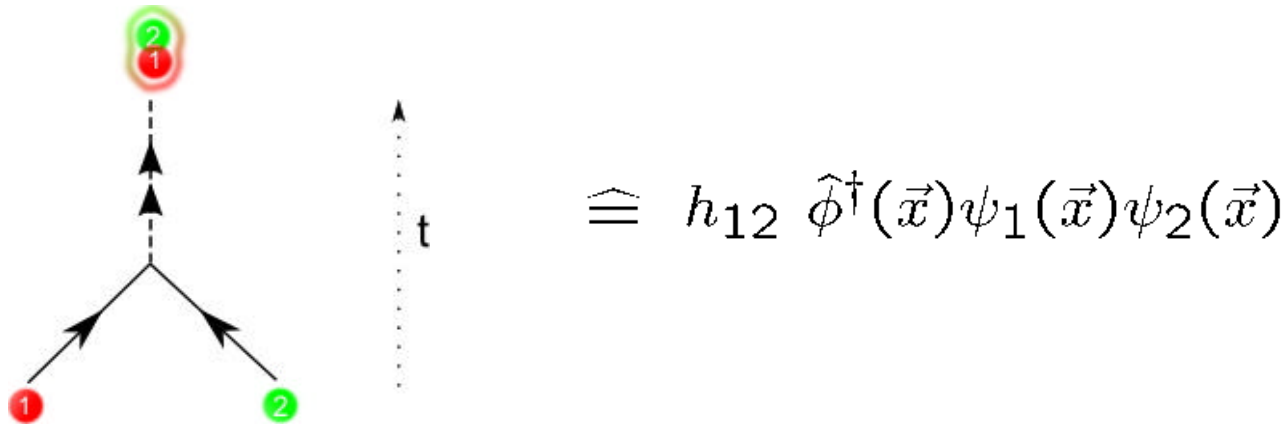


What have we done? We introduced new field(operators) $\hat{\phi}(\vec{x})$, $\hat{\phi}^\dagger(\vec{x})$ describing the creation or annihilation of a bound state of 2 fermions (molecule).

To describe the process of boson formation we need an interaction term

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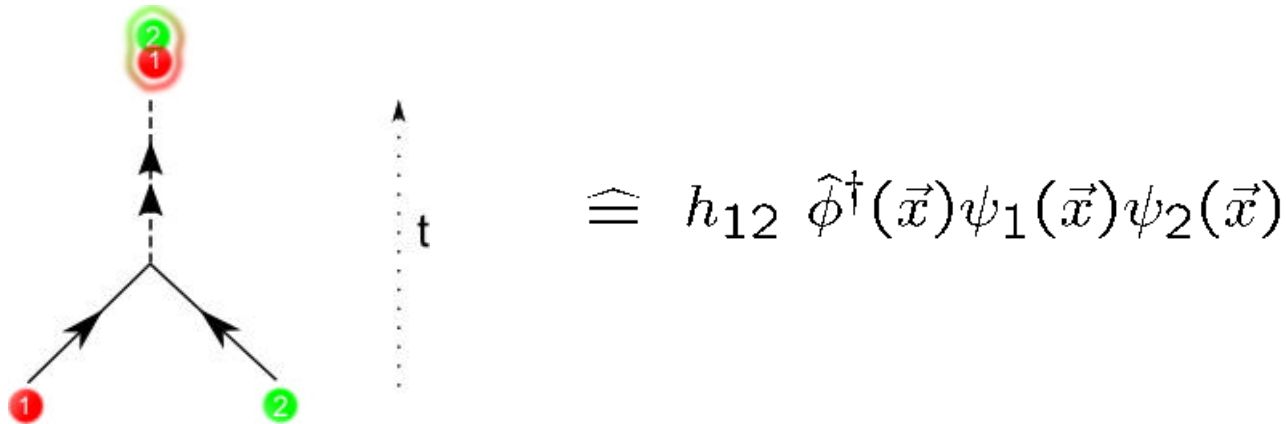
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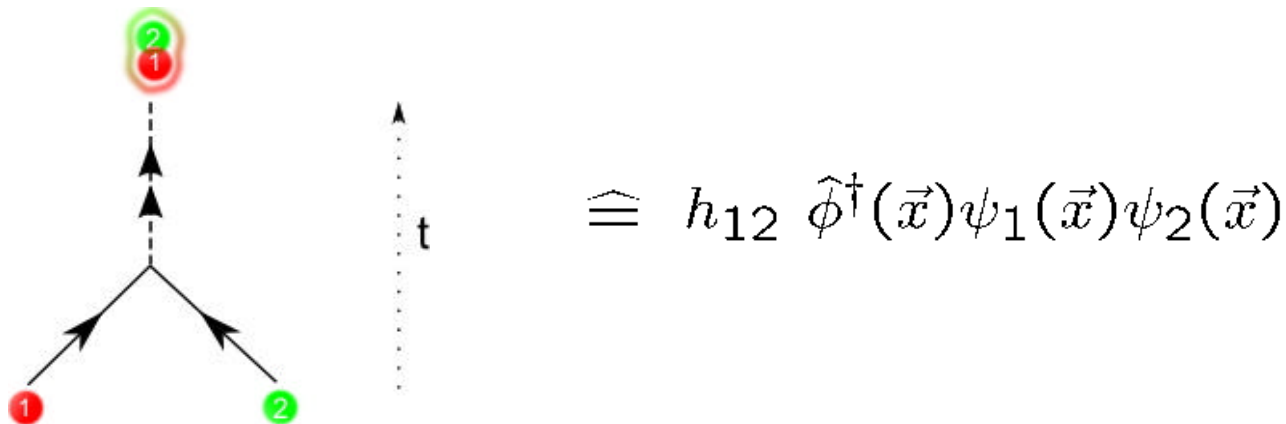
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$$\hat{\psi}_1^\dagger(\vec{x})\hat{\psi}_2^\dagger(\vec{x}) \rightarrow \hat{\phi}^\dagger(\vec{x})$$

$$\lambda_{12} \hat{\psi}_1^\dagger(\vec{x})\hat{\psi}_2^\dagger(\vec{x})\hat{\psi}_2(\vec{x})\hat{\psi}_1(\vec{x}) \rightarrow h_{12} \hat{\phi}^\dagger(\vec{x})\psi_1(\vec{x})\psi_2(\vec{x})$$

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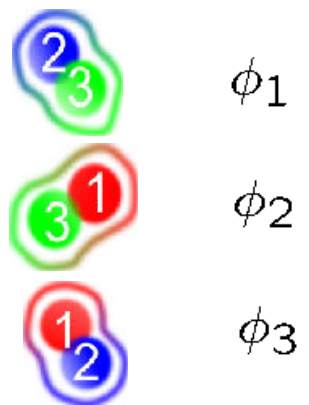
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Remark: This procedure is very much similiar to what you find in mean field BCS theory. Here you perform the replacement

$$\lambda_{12} \hat{\psi}_1^\dagger(\vec{x})\hat{\psi}_2^\dagger(\vec{x})\hat{\psi}_2(\vec{x})\hat{\psi}_1(\vec{x}) \rightarrow \Delta_{12}^* \hat{\psi}_2(\vec{x})\hat{\psi}_1(\vec{x})$$

$$\Delta_{12}^* = \lambda \sum_{\vec{k}} \langle \hat{\psi}_{1,\vec{k}}^\dagger \hat{\psi}_{2,-\vec{k}}^\dagger \rangle$$

Hold on! We missed something. There are three possibilities to form bosonic molecules:



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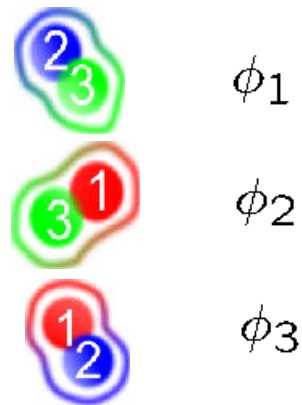
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Ending up with a Hamiltonian of the type

$$\begin{aligned}
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 & + \int d^3x h_{12} \left(\hat{\phi}_3^{\dagger}(\vec{x}) \psi_1(\vec{x}) \psi_2(\vec{x}) + h.c. \right) \\
 & - \int d^3x h_{13} \left(\hat{\phi}_2^{\dagger}(\vec{x}) \psi_1(\vec{x}) \psi_3(\vec{x}) + h.c. \right) \\
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 \end{aligned}$$

Question: *Which bosons will develop?*

Each boson field $\hat{\phi}_i$ can develop an expectation value, giving the corresponding particle density:

$$n_i = \langle 0 | \hat{\phi}_i^* \hat{\phi}_i | 0 \rangle$$

Answer: Statistical physics – Minimize the free energy with respect to this expectation value

$$F = -k_B T \ln(\text{Tr}[exp^{-\beta(\hat{H} - \mu\hat{N})}])$$

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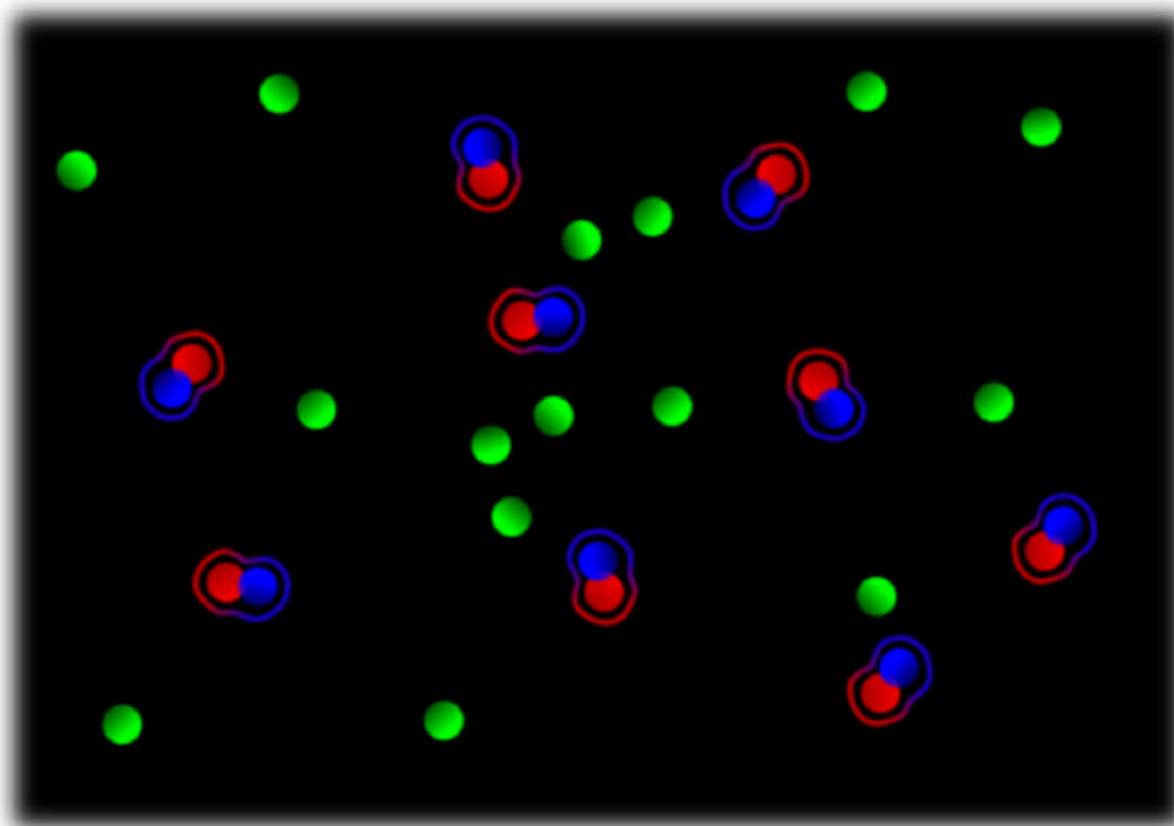
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Depending on the **masses** of the fermions, the **scattering lengths** (couplings), the **densities** or a given **trap geometry**, it might be energetically favorable to:

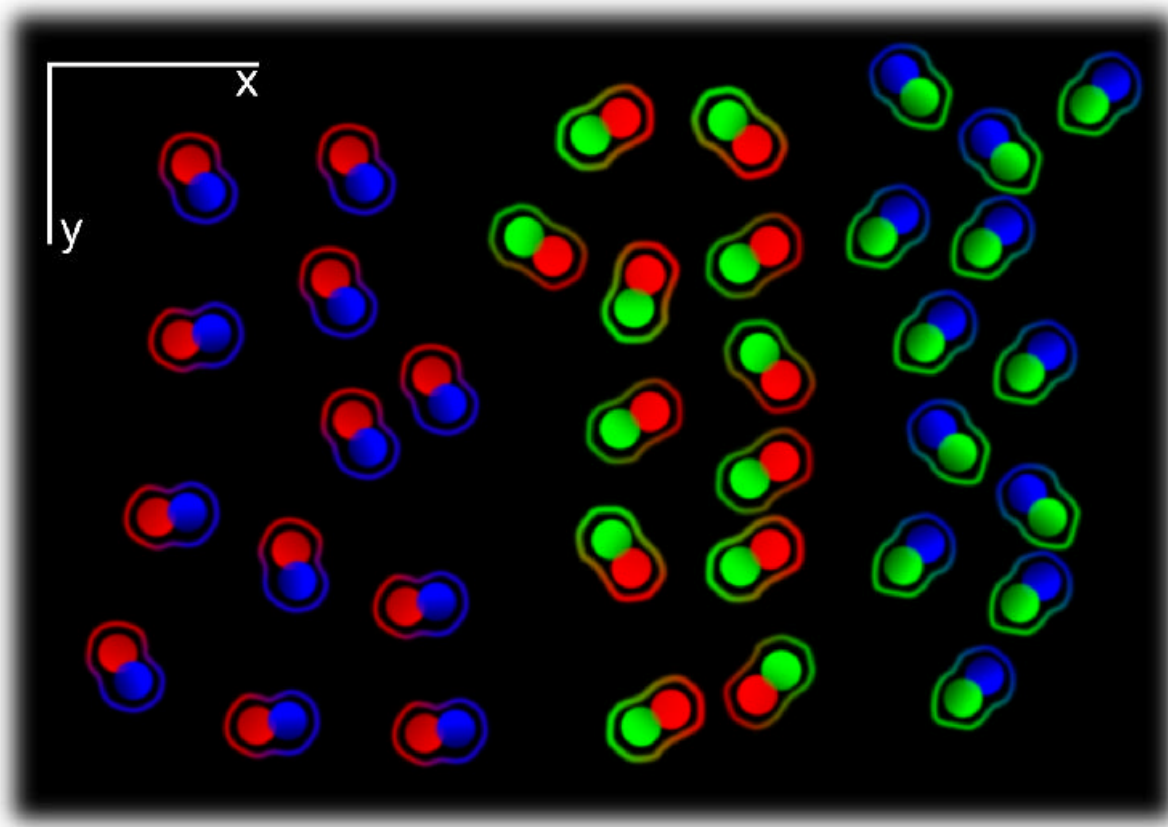
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- form bosons of type 1, 2 or 3 leaving the other one unbound



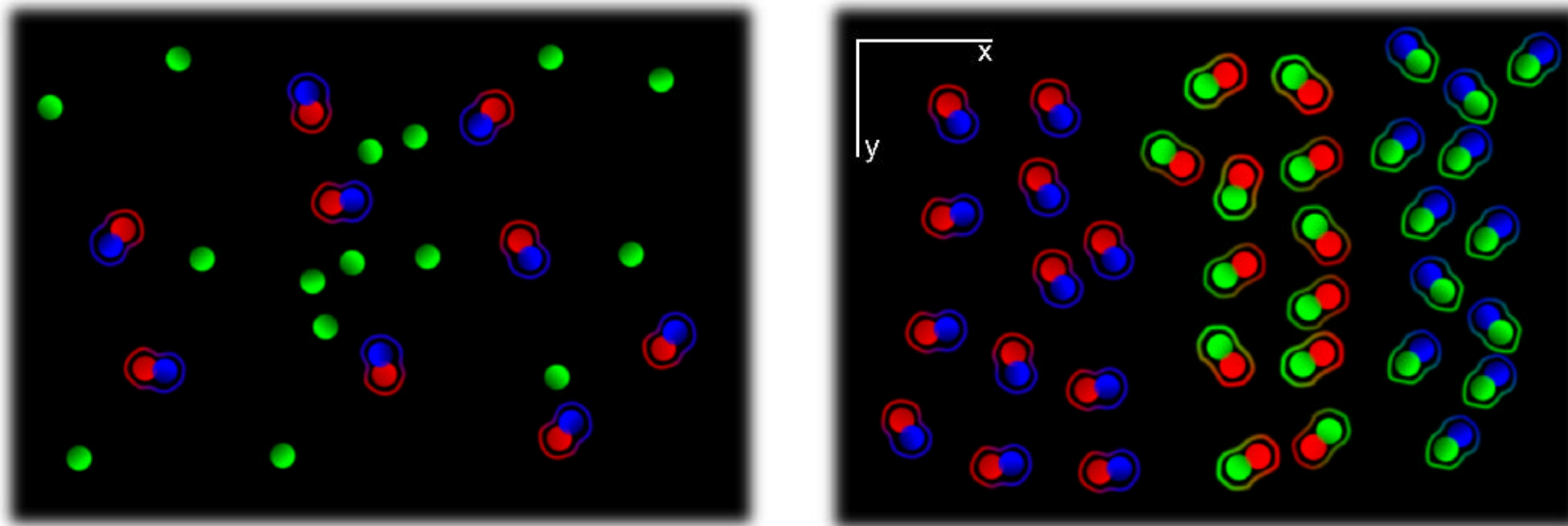
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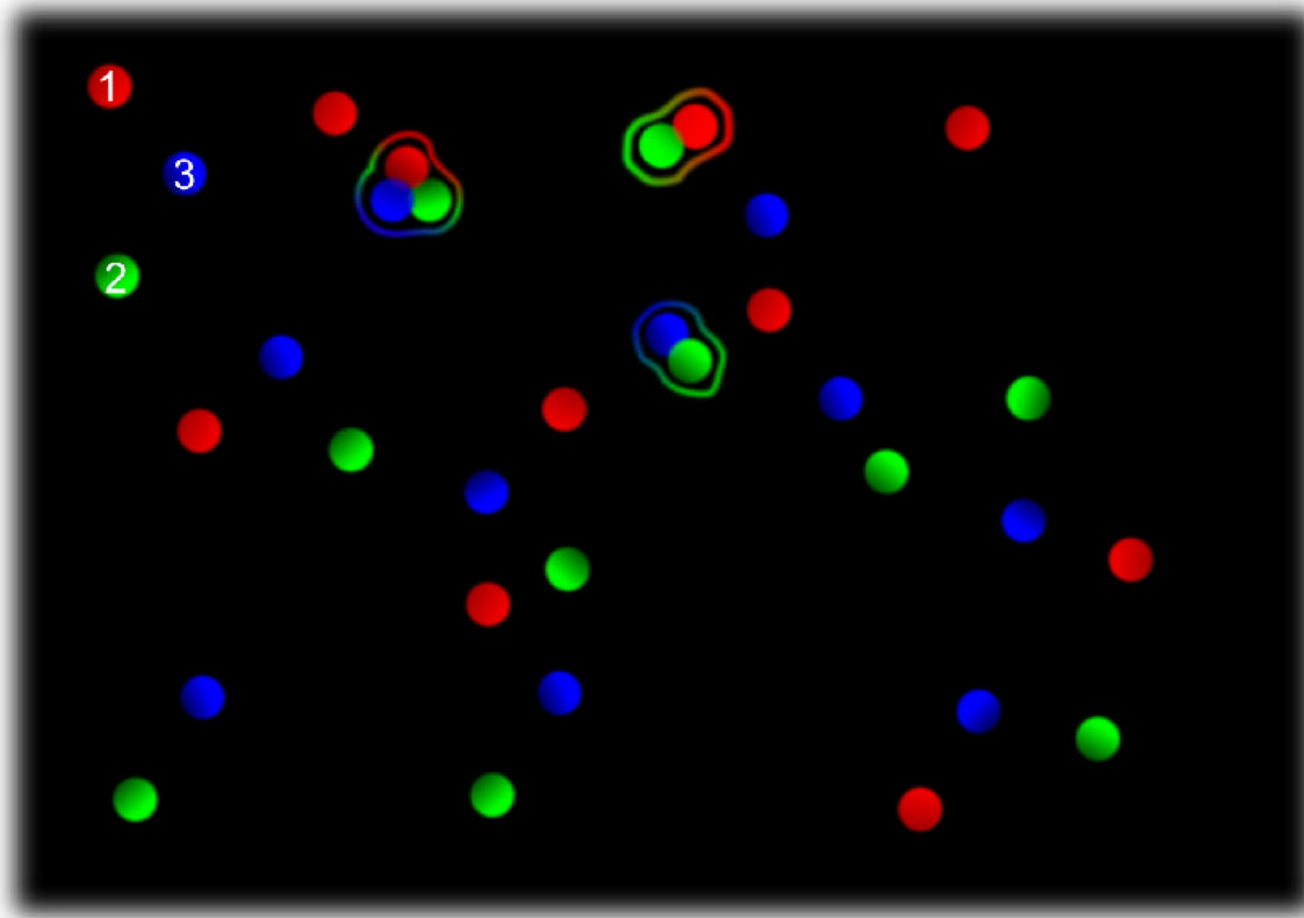
- If all bosons are energetically equivalent, nature will randomly decide for one bosonic type: *Spontaneous* symmetry breaking (SSB).

Important:

These bosons are the ones responsible for developing a BEC or BCS phase!

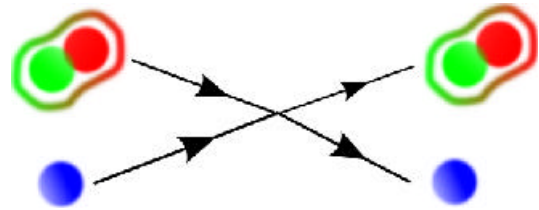
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What's about the promised *trion*?

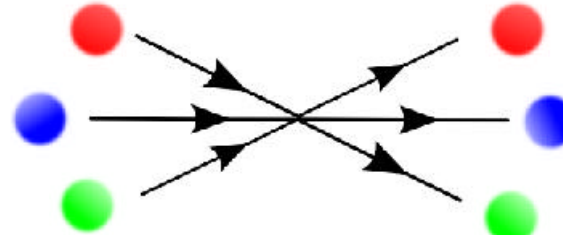


What's about the promised *trion*?

In our Hamiltonian we missed terms (allowed by quantum numbers) describing the scattering processes



$$\propto \hat{\phi}_3^\dagger \hat{\psi}_3^\dagger \hat{\phi}_3 \hat{\psi}_3$$

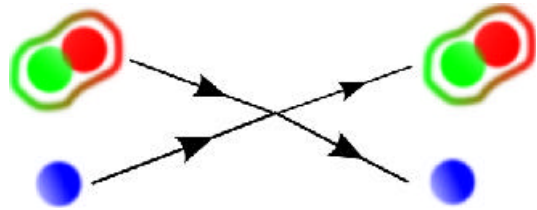


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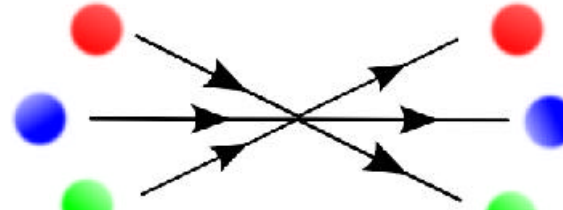
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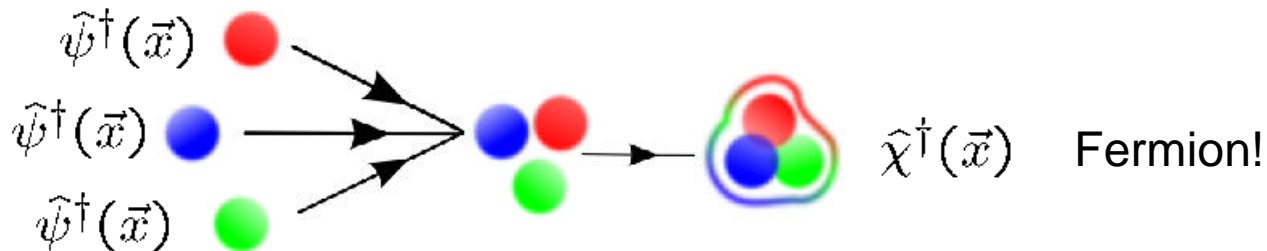
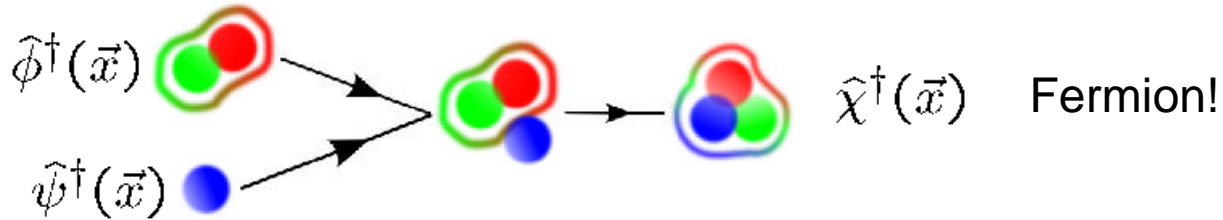


$$\propto \hat{\phi}_3^\dagger \hat{\psi}_3^\dagger \hat{\phi}_3 \hat{\psi}_3$$



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Analogous to before we can do replacements...



The Question:

Will it be energetically favorable for the system to build these Trions?

The Hamiltonian under final consideration is

$$\begin{aligned}
 \hat{H} = & \sum_{\sigma=1,2,3} \int d^3x \hat{\psi}_\sigma^\dagger(\vec{x}) \left(-\frac{\hbar^2}{2m_\sigma} \nabla^2 + \Delta_\phi \right) \hat{\psi}_\sigma(\vec{x}) + \sum_{\sigma=1,2,3} \int d^3x \hat{\phi}_\sigma^\dagger(\vec{x}) \left(-\frac{\hbar^2}{4m} \nabla^2 + \Delta_\phi \right) \hat{\phi}_\sigma(\vec{x}) \\
 & + \int d^3x \hat{\chi}^\dagger(\vec{x}) \left(-\frac{\hbar^2}{6m} \nabla^2 + \Delta_\chi \right) \hat{\chi}(\vec{x}) \\
 & + \sum_{perm.} \int d^3x h_{12} (\hat{\phi}_3^\dagger(\vec{x}) \hat{\psi}_1(\vec{x}) \hat{\psi}_2(\vec{x}) + h.c.) + \sum_{\sigma=1,2,3} \int d^3x g_{\sigma\sigma} (\hat{\chi}^\dagger(\vec{x}) \hat{\phi}_\sigma(\vec{x}) \hat{\psi}_\sigma(\vec{x}) + h.c.)
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$$H \xrightarrow{\text{Legendretransf.}} \mathcal{L} \xrightarrow{\text{Action}} S = \int \mathcal{L} \xrightarrow{\text{Path integral for Quantum Statistics}} \int \mathcal{D}e^{-S} \rightarrow F$$

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$$\begin{aligned} \hat{H} = & \sum_{\sigma=1,2,3} \int d^3x \hat{\psi}_\sigma^\dagger(\vec{x}) \left(-\frac{\hbar^2}{2m_\sigma} \nabla^2 + \Delta_\phi \right) \hat{\psi}_\sigma(\vec{x}) + \sum_{\sigma=1,2,3} \int d^3x \hat{\phi}_\sigma^\dagger(\vec{x}) \left(-\frac{\hbar^2}{4m} \nabla^2 + \Delta_\phi \right) \hat{\phi}_\sigma(\vec{x}) \\ & + \int d^3x \hat{\chi}^\dagger(\vec{x}) \left(-\frac{\hbar^2}{6m} \nabla^2 + \Delta_\chi \right) \hat{\chi}(\vec{x}) \\ & + \sum_{perm.} \int d^3x h_{12} (\hat{\phi}_3^\dagger(\vec{x}) \hat{\psi}_1(\vec{x}) \hat{\psi}_2(\vec{x}) + h.c.) + \sum_{\sigma=1,2,3} \int d^3x g_{\sigma\sigma} (\hat{\chi}^\dagger(\vec{x}) \hat{\phi}_\sigma(\vec{x}) \hat{\psi}_\sigma(\vec{x}) + h.c.) \end{aligned}$$

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Free energy as path integral really hard to compute. Other method:

Functional renormalization group (FRG)

(gives differential form of path integral)

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$$+ \int d^3x \hat{\chi}^\dagger(\vec{x}) \left(-\frac{\hbar^2}{6m} \nabla^2 + \Delta_\chi \right) \hat{\chi}(\vec{x})$$

$$+ \sum_{perm.} \int d^3x h_{12} (\hat{\phi}_3^\dagger(\vec{x}) \hat{\psi}_1(\vec{x}) \hat{\psi}_2(\vec{x}) + h.c.) + \sum_{\sigma=1,2,3} \int d^3x g_{\sigma\sigma} (\hat{\chi}^\dagger(\vec{x}) \hat{\phi}_\sigma(\vec{x}) \hat{\psi}_\sigma(\vec{x}) + h.c.)$$

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The Question:

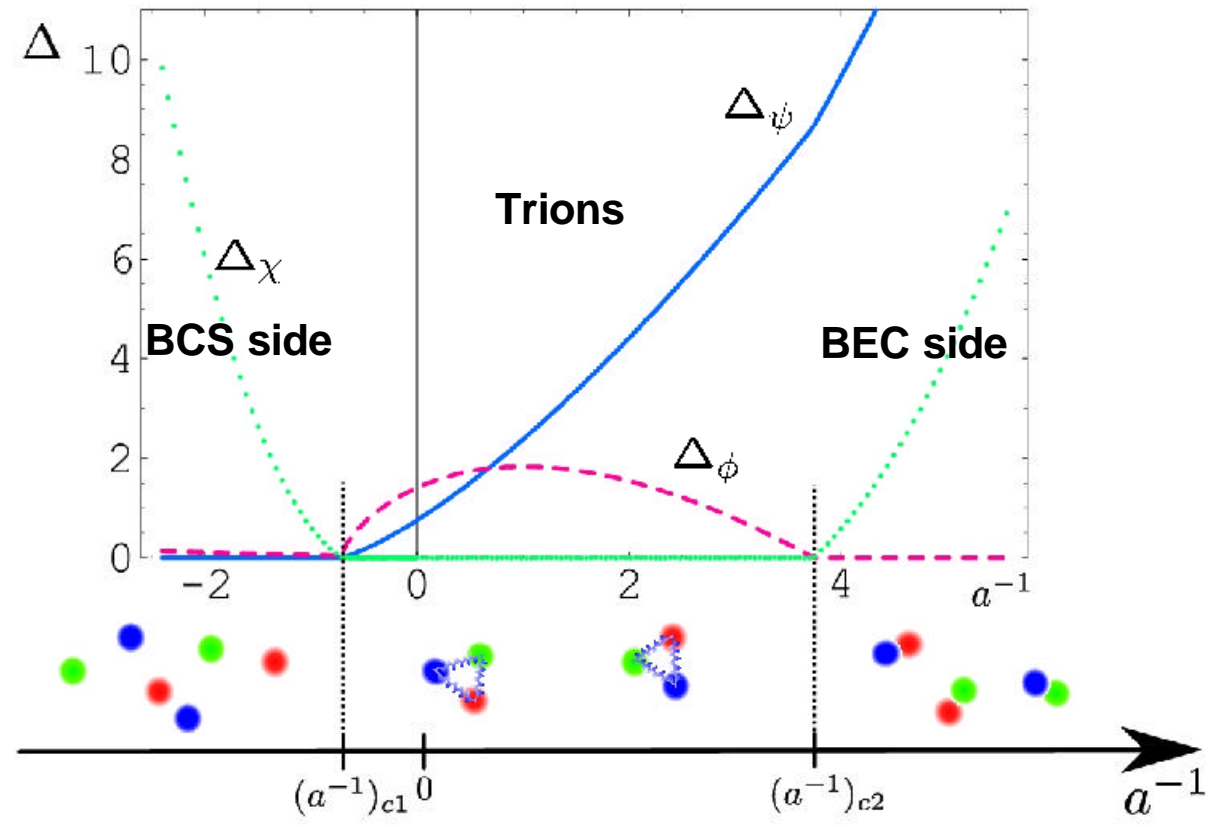
Will it be energetically favorable for the system to build these Trions?

Very preliminary **Answer** for the vacuum:

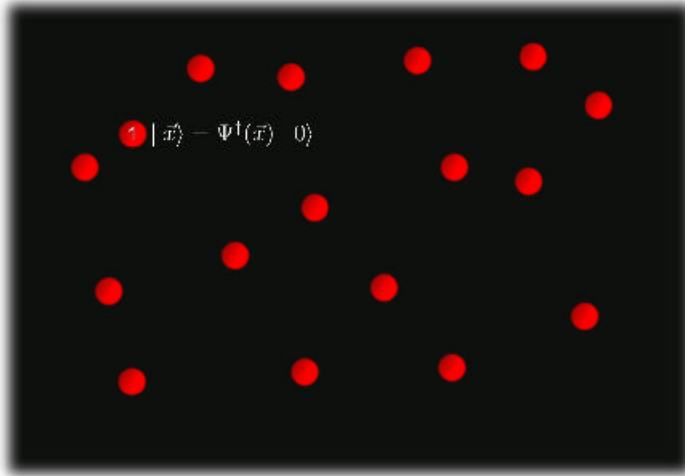
The Question:

Will it be energetically favorable for the system to build these Trions?

Very preliminary **Answer** for the vacuum:



Quite possible!



$\hat{\psi}^\dagger(\vec{x})$ applied to the vacuum creates a new particle at position x .

$\hat{\psi}(\vec{x})$ annihilates a particle at position x .

$|0\rangle$ denotes the vacuum.

$|\vec{x}\rangle$ represents a particle at position x .

How does the Hamiltonian for such a system look like?

$$\hat{H} = \int d^3x \hat{\psi}^\dagger(\vec{x}) \left(-\frac{\hbar^2}{2m} \Delta + V(\vec{x}) \right) \hat{\psi}(\vec{x})$$

But we probably expected something like the Schrödinger Equation

$$i\hbar \partial_t \hat{\psi}(\vec{x}, t) = \left(-\frac{\hbar^2}{2m} \Delta + V(\vec{x}) \right) \hat{\psi}(\vec{x}, t)$$

$\hat{\psi}(\vec{x})$ is an operator. Thus it has to fulfil the Heisenberg equation of motion:

$$i\hbar \partial_t \hat{O} = [\hat{O}, \hat{\psi}]$$

Plugging in $\hat{\psi}(\vec{x})$, yields the Schrödinger Equation for the operator

$$i\hbar \partial_t \hat{\psi}(\vec{x}, t) = \left(-\frac{\hbar^2}{2m} \Delta + V(\vec{x}) \right) \hat{\psi}(\vec{x}, t)$$

Now we should be more comfortable with this (2nd) quantization method.



In the Hamiltonian this interaction corresponds to a term of the form

$$\langle 0 | \hat{\psi}_2(\vec{x}_2) \hat{\psi}_1(\vec{x}_1) [\lambda_{12} \hat{\psi}_1^\dagger(\vec{x}) \hat{\psi}_2^\dagger(\vec{x}) \hat{\psi}_2(\vec{x}) \hat{\psi}_1(\vec{x})] \hat{\psi}_1^\dagger(\vec{x}_1) \hat{\psi}_2^\dagger(\vec{x}_2) | 0 \rangle$$

$\xleftarrow{\neq 0}$
 $\xrightarrow{\neq 0}$

$\hat{\psi}_i(\vec{x})$ describes an incoming particle.

$\hat{\psi}_i^\dagger(\vec{x})$ describes an outgoing particle.

λ_{12} determines the strength of interaction (\rightarrow scattering length).

Remark:

$$\hat{H} = \sum_{\sigma=1,2,3} \int d^3x \hat{\psi}_\sigma^\dagger(\vec{x}) \left(-\frac{\hbar^2}{2m_\sigma} \Delta \right) \hat{\psi}_\sigma(\vec{x})$$

$$+ \int d^3x \lambda_{12} \hat{\psi}_1^\dagger(\vec{x}) \hat{\psi}_2^\dagger(\vec{x}) \hat{\psi}_2(\vec{x}) \hat{\psi}_1(\vec{x})$$

$$+ \int d^3x \lambda_{13} \hat{\psi}_1^\dagger(\vec{x}) \hat{\psi}_3^\dagger(\vec{x}) \hat{\psi}_3(\vec{x}) \hat{\psi}_1(\vec{x})$$

$$+ \int d^3x \lambda_{23} \hat{\psi}_2^\dagger(\vec{x}) \hat{\psi}_3^\dagger(\vec{x}) \hat{\psi}_3(\vec{x}) \hat{\psi}_2(\vec{x})$$

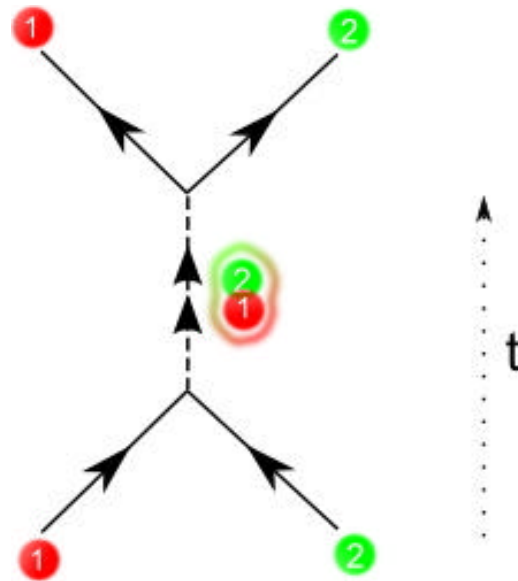
$$\hat{H} = \sum_{\sigma=1,2,3} \int d^3x \hat{\psi}_\sigma^\dagger(\vec{x}) \left(-\frac{\hbar^2}{2m_\sigma} \Delta \right) \hat{\psi}_\sigma(\vec{x})$$

$$+ \int d^3x h_{12} (\hat{\phi}_3^\dagger(\vec{x}) \psi_1(\vec{x}) \psi_2(\vec{x}) + h.c.)$$

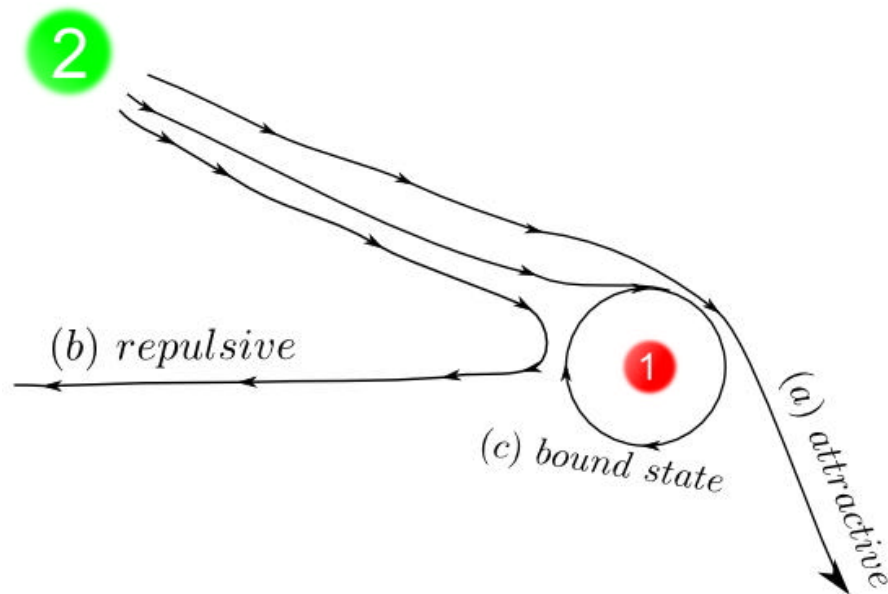
$$- \int d^3x h_{13} (\hat{\phi}_2^\dagger(\vec{x}) \psi_1(\vec{x}) \psi_3(\vec{x}) + h.c.)$$

$$+ \int d^3x h_{23} (\hat{\phi}_1^\dagger(\vec{x}) \psi_2(\vec{x}) \psi_3(\vec{x}) + h.c.)$$

Where is the 4 fermion interaction? Answer:



Recall



This time we are interested in process (c), describing molecule formation.
 In diagrammatic language this translates to