

Universität Heidelberg

"Measurement of electrons from beauty hadron decays at $\sqrt{s} = 7$ TeV and $\sqrt{s} = 2.76$ TeV" **Summary of**

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January 28, 2013 HD group meeting **Motivation of the measurement**

Analysis method and key components of this analysis

Results

pT differential cross section and comparison to the alternative methods

Total cross section

Motivation of $b \rightarrow e$ measurement in pp collisions

- Heavy quarks are produced dominantly via initial hard parton scatterings, therefore their production cross section constitute a prime benchmark for the pQCD calculations in a new energy domain
- In pp collisions at LHC, investigated at \sqrt{s} = 7 TeV in various channels
	- Beauty hadrons at forward rapidity with LHCb and at mid rapidity with CMS (high pt only)
	- Electrons and muons from semi-leptonic decays of heavy flavour hadrons with ATLAS (high p_t only) and with ALICE (down to low p_t)
	- At low p_t , J/ ψ from beauty hadron decays at mid rapidity with ALICE
	- At low p_t , D mesons at mid rapidity with ALICE
	- Good agreement with higher oder pQCD calculation

✓What is missing...

 \blacktriangleright Separation of leptons from charm and beauty hadron decays at low p_t : important for the total beauty production cross section and provide baseline for the PbPb measurement

Papers

measured at mid-rapidity ($|y| < 0.8$) in the transverse momentum range $1 < p_T < 8$ GeV/*c* with the ALICE experiment at the CERN LHC in pp collisions at a center of mass energy \sqrt{s} = 7 TeV using an integrated luminosity of 2.2 nb⁻¹. Electrons from beauty hadron decays were selected based on the displacement of the decay vertex from the collision vertex. A perturbative QCD calculation agrees with the measurement within uncertainties. The data were extrapolated to the full phase space to determine the total cross section for the production of beauty quark-antiquark pairs.

Measurement of electrons from beauty hadron decays in pp.

collisions at $\sqrt{s} = 7$ TeV

The ALICE collaboration¹

⁵ *Keywords:* LHC, ALICE experiment, pp collisions, Single electrons, Heavy flavour ⁶ production, Beauty production

⁷ The measurement of heavy-flavor (charm and beauty) production in proton–proton (pp) ⁸ collisions at the CERN Large Hadron Collider (LHC) provides a crucial testing ground for quantum chromodynamics (QCD), the theory of strong interactions, in a new high-¹⁰ energy regime. Because of their large masses heavy quarks are mainly produced via initial hard parton-parton collisions, even at low transverse momenta p_T . Therefore, heavy-flavor ¹² production cross sections constitute a prime benchmark for perturbative QCD (pQCD) calculations. Furthermore, heavy-flavor measurements in pp collisions provide a mandatory ¹⁴ baseline for corresponding studies in nucleus-nucleus collisions. Heavy quark observables are sensitive to the properties of the strongly interacting partonic medium which is produced in such collisions.

18 Tevation in boing done with referee final iteration is being done with referee first draft is ready for IRC

The same trend was observed in pp collisions at √ ²¹ *s* = 0*.*2 TeV at RHIC [5, 6].

Beauty production in pp collisions at \sqrt{s} = 2.76 TeV, measured using semi-electronic decays

The ALICE Collaboration

Abstract

The production cross section of electrons from semi-leptonic decays of beauty hadrons has been measured in pp collisions at \sqrt{s} = 2.76 TeV at mid-rapidity (|*y*| < 0.8) and in the transverse momentum range 1-10 GeV/*c* with the ALICE experiment at the CERN LHC. Using an impact parameter analysis based on decay vertices which are displaced from the primary vertex of the collisions, electrons from the decay of beauty hadrons are selected. The production cross section of beauty decay electrons was compared to the result obtained utilizing an alternative method which uses azimuthal correlations of heavy-flavour decay electrons and charged hadrons. We also compare the relative beauty fraction of the total heavy-flavour electron spectrum measured using the correlation technique to that obtained with the impact parameter analysis. In addition, we compare to pQCD predictions in the FONLL framework and the calculation is in agreement within the uncertainties. The result was extrapolated to the full phase space to determine the total $b\bar{b}$ production cross section.

Keywords: LHC, ALICE experiment, pp collisions, Single electrons, Heavy-flavour production, Beauty production

1. Introduction

Heavy-flavour quarks are of particular interest in pp collisions because they are, unlike their lighter counterparts, produced through the initial hard parton-parton scatterings. Therefore, the measurement of their production provides essential tests of perturbative Quantum Chromodynamic (pQCD) calculations. Additionally, these measurements in pp collisions provide the necessary baseline for the equivalent measurements performed in heavy ion collisions. In pp collisions at [√]*^s* ⁼ 2.76 TeV ALICE has mea-

each measurement being in good agreement with pQCD

from semi-electronic decay of beauty hadrons measured in the mid-rapidity region $(|y| < 0.8)$ with the AL-ICE experiment in the range $1 < p_T < 10$ GeV/*c* in pp collisions at \sqrt{s} = 2.76 TeV and present the total bb production cross section based on the extrapolation to full phase space from the measured p_T -differential cross section. The results are compared to the predictions from FONLL pQCD corresponding calculations [5]. The results are measured primarily using an impact parameter analysis which takes advantage of the relatively long lifetime of beauty hadrons com-

flavour electron spectrum to compute the production spectrum to compute the production of production and production of production of the production of

[1], heavy-flavour hadrons via semi-leptonic decays to electrons (mid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-rapid-ra \overline{a} and charged of \overline{b} $\bm{\mathsf{A}}$ dons ($\bm{\mathsf{A}}$ to extract the relative beauty contri- \mathbf{f}

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Agreement between data and MC of key variable(impact parameter) (1)

Use of conversion electrons identified by V0 finder - Only identifiable source

Agreement between data and MC of key variable(impact parameter) (2)

- *p*^T range 1*.*5 *< p*^T *<* 2 GeV/*c* (left) and 2 *< p*^T *<* 2*.*5 GeV/*c* (right) (*d*⁰ was multiplied with the charge of \bullet The d₀ distribution of the data sample is well described by the cocktail of signal and backgrounds (The differences between the data and the cocktail of the signal and backgrounds are consistent with statistical variations). ⁴³ The systematic uncertainties are 30-40% at 1-1.5 GeV/*c*. In addition, we have statistical uncertainties of
	- The yields of signal and backgrounds obtained by this procedure agree with those obtained in this analysis within statistical uncertainties. *y* and procodure agree with these ⁴⁷ from data. This corresponds to a variation of the signal by about *±*20% (described in lines 172-176).

Agreement between data and MC of key variable(impact parameter) (4)

Raw spectra using the cut based on the impact parameter significance

• At 2<pt<6, both gives **similar good agreement between data and MC (within statistical fluctuation)**

1st pixel

from primary pions I EXKPS **vertex outside of the 1st pixel?**

• Electron production vertex distribution with different hit requirement on pixel (standard HFE cuts applied)

• **"Hits on both pixel" requirement is effective to reduce the fake tracks produced outside of the first pixel.**

Impact parameter dist. with different pixel hit requirements

• **Better to remove them as much as possible to reduce any possible discrepancy of fakes between data and MC, and to increase S/B**

pt

pt

Hadron contamination: less suffer than inclusive analysis

PID with TPC-TOF(|σ TOF|<3, 0(-1)<σ TPC<3)

hadrons and from photon conversions obtained from PYTHIA simulations in the electron *p*^T range 1 *< p*^T *<* The PYTHIA simulation does not reproduce precisely the p_T-differential yields of background sources measured in data. Therefore, the p_T distributions of the relevant electron sources in PYTHIA were re-weighted to match the distributions measured with ALICE

The production cross sections of π^0 and η mesons, the dominant sources of electrons from Dalitz decays and from photons which convert in material into e⁺e[−] pairs, were measured with ALICE in pp collisions

 \bullet D^0 , D^+ , and D^+ _s meson production cross sections were measured with ALICE in $1 < pT < 16$ GeV/c, $1 < pT < 24$ GeV/c, and $2 < pT < 12$ GeV/c, respectively.

• Based on a FONLL pQCD calculation the measured p_Tdifferential cross sections were extrapolated to $pT = 50$ GeV/c.

- ² • Contribution from the unmeasured high- p_T region to the ≤ 10% for electrons with $p_T < 8$ GeV/c.
- $\frac{1}{10}$ $\frac{1}{10}$ $\overline{\text{max}}$ and $\overline{\text{max}}$ in the same ratio $\sigma(\Lambda_c)/\sigma(D^0 + D^+)$ from ZEUS. • A contribution from Λ_c decays was included using a
- $\begin{bmatrix} \begin{matrix} 1 & 1 & 1 \\ 1 & 1 & 1 \end{matrix} \end{bmatrix}$ For 2.76 TeV, D meson cross sections were obtained by applying 10³ december 113 **d** based and different sources are estimated using the different sources are estimated using the measured parameter $\frac{10^4}{2}$

 $\frac{1}{2}$ a √s -scaling to the cross sections measured at √s = 7 TeV.

 $\frac{2}{3}$ $\frac{1}{2}$ $\frac{1}{2}$ **E** a MC scaled 7 TeV spectra using the scaled 7 TeV $\frac{1}{1+\frac{1}{2}+\frac{1}{3}+\cdots+\frac{1}{4}+\cdots+\frac{1}{5}+\cdots+\frac{1}{5}+\cdots+\cdots+\frac{1}{5}+\cdots}}$ measurements within statistical uncertainties.

Variation of the selection cuts to estimate systematic uncertainties particle in the selection car

NonHFE backgrounds: Uncertainty of \mathbb{R}^{10} input meson spectra (same spectra as solid in Table 5.10 used for inclusive HFE analysis) $\frac{1}{2}$ and $\frac{1}{2}$ 267.7 and charm hadron decays) was estimated as the section 3.8 . 2688 For each variation of the selection criteria, the background subtracted electron spectrum is fully corrected electron spectrum is fully corrected electron spectrum is fully corrected electron spectrum is fully corr Charm backgrounds: Uncertainty **with the resulting spectra are compared by inspecting and the compared by inspectra and the compared** of D meson spectra 10^{-6} $\begin{bmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$ 27777 spectra. A general systematic uncertainty of 2% 2% 2% **) 2 (mb/(GeV/c) |y|<0.5 dy T /dp** m **2 d T p**/ **1/2 -7 10 0.5 10⁻⁶** 10^{-5} 10^{-4} **-3 10 -2 10** \sum 10⁻¹ ALICE $D \rightarrow e$ \rightarrow FONLL D \rightarrow e $pp, \sqrt{s} = 7$ TeV **Data/FONLL -6 10** pp, s = 7 TeV **0** 2 4 6 8 10
 p_t (GeV/c) **Formal 0**

Overview over the contributions to the systematic uncertainties

MinJung, University of Heidelberg January 28th, 2013 ²¹³ BR^Hb→^e + BR^Hb→Hc→^e = 0*.*205 *±* 0*.*007 [25]. The related uncertainty was obtained as the

Alternative measurements

p
Pecuon agree with the result obtained using the impact parameter cut. Within the uncertainties, these alternative approaches to access the beauty e lectron p_T -differential cross section agree

) <u>ณ</u>

dy) (mb/(GeV/c)

1/(2πp) d² o/(dp _τ

The azimuthal correlation between heavy-flavour decays and correlation between heavy-flavour decays and the azimuthal correlation between heavy-flavour decays and the and charged the assemble of the and charged the state

p_T-differential cross sections and weighted average AINNTAN AVARANA Electrons from beauty hadron decays in pp collisions at √*s* = 2*.*76 TeV 21

- \bullet Over the full accesible nt range \bullet ONI Accible nt range FONIL predictions are in good agreement • Over the full accessible pt range, FONLL predictions are in good agreement with the data
- lines indicate the corresponding FONLL predictions (uncertainties) [20]. Ratios of the data and the FONLL calculations are shown in the Gastrons of the Beauty hadron decays take over from c the dashed lines indicate the FONLL uncertainties. (d) Measured ratio of electrons from beauty and charm hoquu flouar bodron • Beauty hadron decays take over from charm as the dominant source of electrons from \sim 1*H* \sim ⁴⁰² *7 TeV [7])* heavy-flavor hadron decays close to electron p_T of 4 GeV/c. ³⁷⁹ In Fig. 24, the measured *p*T-differential cross section of electrons from beauty hadron decays (black

Calculation of total cross section charm charm of the mid-⁵³⁶ Using eq. 3 with the branching ratio of charm into electrons and the kinematical correction factor men- 552

- \bullet p_t range of measured p_t differential cross section:
- \bullet 1(0.5) < p_t <8 GeV/c for b \rightarrow e (b,c \rightarrow e)

a Extrapolation begand an EQNAU obene
- = 889*.*18*±*36*.*85(*stat*) [−]339*.*16(*sys*) [−]127*.*79(*extr.*)*±*37*.*35(*norm*)*±*36*.*74(*br.*)*µ^b* • Extrapolation based on FONNL shape
	- \bullet down to 0 p $_{\rm t}$ for d σ (b \rightarrow e)/dy (51% unmeasured based o *dy* ∆*y* • down to 0 p_t for d σ (b \rightarrow e)/dy (51% unmeasured based on FONLL)
	- \bullet to full rapidity for total bp cross section f<mark>or total</mark> l +43 +24 [−]46(*sys*) \bullet to full rapidity for total bb cross section

$$
\frac{d\sigma(b,c \to e)}{dy} = \frac{\sigma_{vis} * \beta_{b,c}}{\Delta y} \qquad \text{where} \qquad \beta_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 \text{GeV}/c < p_t < \infty, y_{min} < y < y_{max})}{\sigma_{vis, \text{FONLL}}} \n\frac{d\sigma_{b,c}}{dy} = \frac{\sigma_{vis} * \gamma_{b,c}}{\Delta y} \qquad \text{where} \qquad \gamma_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c, 0 \text{GeV}/c < p_t < \infty, y_{min} < y < y_{max})}{\sigma_{vis, \text{FONLL}}} \n\sigma_{b,c} = \frac{\alpha_{b,c} * \sigma_{vis}}{BR(b,c \to e)} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c
$$

- Statistical, Systematic and normalization uncertainties: \overline{a} **d** *d e*)
	- scaled by multiplication factors sections section obtained after variation obtained after variation obtained after variation is compared after variation obtained after variation is compared with a section of the mid-
	-
	- Extrapolation uncertainties:
• each peremeter is resoleulated for different veriction of mass, sools and PDE in FONLL • each parameter is recalculated for different variation of mass, scale and PDF in FONLL [−]0*.*40(*extr*)*±*0*.*29(*norm*)*µ^b* $\bf 5$ recalculated for different variation of mass, scale and $\bf P$ DF in F **or different variation of mass, scale**
- cross section obtained after variation is compared to the cross section obtained using the central value from FONLL and quadratically summed $\frac{1}{2}$ electrons from FONIII and quadratically cummed ⁵⁵⁰ gible compared to the charm contribution for 0*.*7*GeV/c < p*^t *<* 1*GeV/c*. One obtains with this method ⁵⁵⁸ and for the corresponding cross section of electrons from charm hadron decays *d*σ*b,^c* $\overline{}$ 19 *different variation of mass, scale and PDI+111 ONLL
tion is compared to the cross section obtained using*

Total cross sections (7 TeV) σ*vis,FONLL* ⁵⁵⁶ For charm the y-range is *|y| <* 0*.*5 and for beauty it is *|y| <* 0*.*8. One obtains for the cross section of prediction. The uncertainty obtained from this is added linearly to the extrapolation error. The total *bb*¯ ⁵²⁴ Total cross sections (7 TeV) ²²⁴ dictions are in good agreement with the data. At low τ due to the subtraction of τ due to the subtraction of light hadron decay background τ 101al cross sections (7 TeV)

Beauty $\overline{1}$ Beauty for $\overline{1}$ $\overline{2}$ $\overline{3}$ $\overline{4}$ $\overline{2}$ $\overline{4}$ $\overline{2}$ $\overline{3}$ $\overline{4}$ $\overline{2}$ $\overline{4}$ $\overline{2}$ $\overline{2}$ $\overline{3}$ $\overline{4}$ $\overline{2}$ $\overline{3}$ $\overline{2}$ $\overline{3}$ $\overline{4}$ $\overline{2}$ $\overline{3}$ \overline

 $\sigma_{b\rightarrow e}(p_t > 1GeV/c, |y| < 0.8) = 6.61 \pm 0.54(stat)_{-1.86}^{+1.92}(sys) \pm 0.231(norm)\mu b$ ⁵⁵⁷ electrons from beauty hadron decays at midrapidity γ*b,^c* = $\frac{1}{2}$ ($p_t > 1$ GeV / c, $|y| < 0.8$) = 6.61 \pm 0.54(stat) $^{+1.92}_{-1.86}$ (sys) \pm 0.231(norm) μb <u>beauty</u>
225 *p*t, electrons from heavy-flavor hadron decays original telectron decays original telectron decays original telectron decays of the communication of the communication of the communication of the communication $\Omega_{b\to e}(p_t > 1$ GeV $/c, |y| < 0.8$) = 0.01 \pm 0.34(*stdt*) $_{-1.86}$ (*sys*) \pm 0.2 σ_{12} (p ≥ 1 CeV/c |v| ≥ 0.8) = 6.61 + 0.54(stat)^{+1.92}(sys) = $\frac{1}{\sqrt{1-\frac{1}{1-\$ 0.231 (porm) ν $\frac{2}{2}$

$$
\frac{d\sigma(b \to e)}{dy} = 8.40 \pm 0.69(stat)_{-2.36}^{+2.43}(sys)^{+0.19}(extr) \pm 0.29(norm)\mu b
$$

$$
\frac{d\sigma_b}{dy} = 42.25 \pm 3.47(stat)_{-11.88}^{+12.25}(sys)^{+1.08}_{-2.23}(extr) \pm 1.49(norm)\mu b
$$

$$
\sigma_{b\bar{b}} = 280 \pm 23(stat)_{-79}^{+81}(sys)^{+7}_{-8}(extr) \pm 10(BR) \mu b
$$

489 with the same model as it was used for the central prediction of the central prediction. hadron decays measured as 6*.*61*±*0*.*54(stat)+1*.*⁹² [−]1*.*⁸⁶ ²³³ (sys) *^µ*^b ²⁰⁶ charm hadron decays was subtracted from the inclusive ²¹⁹ sistent with the result of a previous measurement of J*/*ψ mesons from beauty hadron decays

dy

 $\sigma_{HF\rightarrow e}(pt > 0.5 \text{GeV} / \text{C}, |y| > 0.5) = 31.13 \pm 3.17 \text{(stat)} + 44 \text{(sys)} + 1.52 \text{(normal)}$ $\sigma_{HF\rightarrow e}(p_t > 0.5 GeV/c, |y| < 0.5) = 37.73 \pm 3.17(stat)_{14.44}^{+13.5}(sys)$ **d**
de $\sigma_{HF\to e}(p_t > 0.5 GeV/c, |y| < 0.5) = 37.73 \pm 3.17(stat)_{14.44}^{+13.3}(sys) \pm 1.32(norm)\mu b$ **Charm**
 $\sigma_{HF\rightarrow e}(p_t > 0.5GeV/c, |y| < 0.5) = 37.73 \pm 3.17(stat)_{14.44}^{+13.3}(sys) \pm 1.32(norm)ub$ $2\pi r \rightarrow e \left(P1 \right)$ $2\pi r \rightarrow e \left(P1 \right)$ $Q_{HF\rightarrow e}(pt > 0.5 \text{GeV } / \text{C}, |y| > 0.5) = 31.13 \pm 3.17 \text{(stat)}$ $(14.44 \text{Vs/s}) \pm 3.5 \text{V}$.

dy Visible cross section of beauty decays are subtracted from visible cross section of charm decay $\frac{200}{200}$ cass section of beauty decays are subtracted from visible cross section of charm decay <u>ty</u>

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7 \text{ TeV.} \int \text{Ldt} = 2.2 \text{ nb}^{-1}
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\n<math display="block</math>

²²⁴ dictions are in good agreement with the data. At low

MinJung, University of Heidelberg **January 28th**, 2013 MinJung, University of Heidelberg **are all 1 and inverse of the covariance matrix** danuary 28th, 2013 $5⁵$ and for the total charm cross section sect

 $\sum w_i = 1$ $\frac{\partial^2}{\partial \theta^2}$ error $(\theta \rightarrow e)$ error $(\theta \rightarrow e)$ error $(\theta \rightarrow f)/\psi$ $\sigma_{error(b\rightarrow e)}^2 \sigma_{error(b\rightarrow J/\psi)}^2$ $\sigma_{error(b\rightarrow J/\psi)}^2$ ⁵⁴⁷ of the visible cross section is 1*GeV/c < p*^t *<* 8*GeV/c*, while for charm it is 0*.*7*GeV/c < p*^t *<* 8*GeV/c*. ⁵⁵⁰ gible compared to the charm contribution for 0*.*7*GeV/c < p*^t *<* 1*GeV/c*. One obtains with this method *b* a single physical quantity *±*
 a single physical quantity *±* Reference: [L.Lyons et al., NIM A 270 \(1988\) 110](http://dx.doi.org/10.1016/0168-9002(88)90018-6) "How to combine correlated estimates of a single physical quantity"

weighted average $\frac{1}{2}$

weighted average of total c

⁵⁶⁶ The total beauty cross section is combined with the one obtained from non-prompt *J/*Ψ [4] via weighted

$$
\bar{\sigma_b} = \sum_{i=0}^1 \sigma_{b,i} w_i
$$

where $\sigma_{b,l}$ are the individual cross sections and w_l are where $\sigma_{b,i}$ are the individual cross sections and w_i are the weighting factors \tilde{A} where $\mathcal{L}_{U,l}$ the ratio intervention of the sections the \mathcal{L}_{l} the

⁵⁶⁶ The total beauty cross section is combined with the one obtained from non-prompt *J/*Ψ [4] via weighted

 $\widehat{\mathbf{e}}$ ↑
ಹ
ಠ 8.9

8.5

8.6

8.7

8.8

Statistical and systematic errors between two measurements are uncorrelated, ⁵⁶⁵ *4.4.3 Weighted average of the total beauty cross section* but extrapolation errors are correlated(via α)

$$
\sigma_{b,c} = \frac{\alpha_{b,c} * \sigma_{vis}}{BR(b,c \to e)} \qquad \text{where} \quad \alpha_{b,c} = \frac{\sigma_{\textit{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{\textit{vis},\textit{FONLL}}}
$$

error matrix

$$
E = \begin{pmatrix} \sigma_{error(b \rightarrow e)}^2 & r\sigma_{error(b \rightarrow e)}\sigma_{error(b \rightarrow J/\psi)} \\ r\sigma_{error(b \rightarrow e)}\sigma_{error(b \rightarrow J/\psi)} & \sigma_{error(b \rightarrow J/\psi)}^2 \end{pmatrix}
$$

Weighted average of total cross section

[−]339*.*16(*sys*) [−]127*.*79(*extr.*)*±*37*.*35(*norm*)*±*36*.*74(*br.*)*µ^b*

Weighted average of total cross section will parameters .

1 find the weights minimizing $\sigma^2 = w^T E w$ subject to $\sum w_i = 1$

$\frac{1}{1}$ weights and errors

$$
w = \frac{E^{-1}u}{u^TE^{-1}u} \qquad \sigma_{\text{stat,sys,extr}}^2 = w^TE_{\text{stat,sys,extr}}w
$$

σ*b,iwi* (11)

 α (b \rightarrow J/ ψ)

4.4 4.45 4.5 4.55 4.6

correlation of α

parameters

 $d\sigma_b/dy$ at mid-rapidity as a function of \sqrt{s} in pp and pp collisions

 $\frac{1}{100}$ Total overal obtained using the contral values $\frac{1}{100}$ ⁴¹⁴ is 1 GeV*/c < p*^t *<* 8 GeV*/c*. One obtains with this method for the total beauty cross section Total cross section for 2.76 TeV (calculated in a same way) ^β*b,^c* ⁼ ^σ*FONLL*(*b,^c* [→] *^e,*0 GeV*/^c < ^p*^t *<* [∞] 2.76 **Te** ⁴¹⁹ Here, the range in rapidity is *|y| <* 0*.*8. One obtains for the cross section of electrons from beauty hadron

 $\sigma_b = 129 \pm 15.2(stat)^{+40.9}_{48.6}(sys)^{+3.38}_{205}(extr) \pm 2.45(norm) \pm 4.36(BR)\mu b$ $\sigma_b = 129 \pm 15.2(stat)_{-48.6}^{+40.9}(sys)_{-3.05}^{+3.38}(extr) \pm 2.45(norm) \pm 4.36(BR)\mu b$

$$
\frac{d\sigma(b\to e)}{dy} = 4.53 \pm 0.53(stat)^{+1.44}_{-1.71}(sys)^{+0.08}_{-0.10}(extr) \pm 0.09(norm)\mu b
$$

within uncertainty

,ymin < y < ymax)

prediction σc¯c = 4*.*76+6*.*⁴⁴ ²⁹² [−]3*.*²⁵ mb [21]. In summary, p_t -differential invariant production cross sections of electrons from beauty and from charm hadron decays were measured in pp collisions at $\sqrt{s} = 7$ TeV. and at 2.76 TeV $\fbox{2}$ The agreement between FONLL predictions and the data reported here suggests that higher order pQCD calculations can reliably describe heavy-flavor production even at low p_t in the highest energy hadron collisions accessible in the laboratory today. Furthermore, these new data provide a crucial baseline for heavy-flavor production studies in the hot and dense matter created in Pb-Pb collisions at the LHC. $\overline{\mathcal{O}}_n$

³⁰⁴ Full author list given at the end of the article.

- Measurement of electrons from beauty hadron decays in pp collisions at $\sqrt{s} = 8$ TeV using TRD Level-1 triggered events
- Measurement of the nuclear modification factor of electrons from beauty hadron decays in Pb-Pb collisions at $\sqrt{sNN} = 2.76$ TeV
- Measurement of the nuclear modification factor of electrons from beauty hadron decays in p-Pb collisions
- Beauty jet identification in pp collisions
- Measurement of beauty jet R_{AA} in Pb-Pb collisions
- Measurement of medium modified fragmentation function of beauty jet in Pb-Pb collisions

BACKUP

$$
\frac{d\sigma(b\rightarrow e)}{dy}
$$

[−]0*.*49(*extr*)*±*0*.*29(*norm*)*µ^b* **(Thanks to Markus Fasel and Ralf)**

Extrapolation down to 0 pt

- FONLL data points are interpolated via TSpline3(Tsallis fit for charm)
- \bullet Fit the measured p_t spectra with the normalization as a free parameter

Statistical uncertainties: same as the one from visible cross section

Systematic uncertainties: one from visible cross section + propagation of systematic uncertainties to extrapolation (sum linearly)

- systematic uncertainties are correlated
- move up and down all data points by the systematic uncertainties
- fit the central FONLL prediction to the data points after moving
- evaluate the difference from the fit of central prediction using central data points

Extrapolation uncertainties:

- fit is done for different variation of mass, scale and PDF in FONLL
- cross section obtained after variation is compared to the cross section obtained using the central value from FONLL and quadratically summed
- fit error from measured data points fit (with statical errors) is quadratically added

$$
\frac{d\sigma_b}{dy} = \frac{1}{BR(b \rightarrow e) * C_{eb}} \frac{d\sigma(b \rightarrow e)}{dy}
$$
\n
$$
BR(b = \text{auty}) = 0.2046 \pm 0.0067
$$
\n
$$
BR(charm) = 0.096 \pm 0.004
$$
\n
$$
\Delta \frac{d\sigma_b}{dy} |_{BR} = \frac{1}{BR^2 * C_{eb}} \frac{d\sigma(b \rightarrow e)}{dy} \Delta(BR)
$$
\n
$$
\Delta \frac{d\sigma_b}{dy} |_{BR} = \frac{1}{BR^2 * C_{eb}} \frac{d\sigma(b \rightarrow e)}{dy} \Delta(BR)
$$

Extapolation based on FONLL

\n
$$
\frac{d\sigma_b}{dy} \quad \text{curve shape}
$$
\n
$$
\sigma_b = \frac{\int \frac{d\sigma_b}{dy} \, dy}{\int_{y_{\text{min}}}^{y_{\text{max}}} \frac{d\sigma_b}{dy} \, dy} \times \left(\frac{d\sigma_b}{dy} \, dx\right) \times \Delta y
$$

⁵⁰⁵ central FONLL prediction. The resulting cross section of electrons from beauty hadron decays per unit Extrapolation error based on line shape of upper and lower limit

$$
\frac{d\sigma(b \to e)}{dy} = 8.24 \pm 0.34(stat)^{+2.17}_{-2.11}(sys)^{+0.33}_{-0.49}(extr) \pm 0.29(norm)\mu b
$$
\n
$$
\frac{d\sigma_b}{dy} = 41.46 \pm 1.71(stat)^{+10.93}_{-10.63}(sys)^{+1.68}_{-2.50}(extr) \pm 1.45(norm) \pm 1.41(BR)\mu b
$$
\n
$$
\sigma_b = 275 \pm 11.3(stat)^{+72.6}_{-70.6}(sys)^{+13.2}_{-21.7}(extr) \pm 9.64(norm) \pm 9.35(BR)\mu b
$$

¹⁴*.*⁴⁴ (*sys*)*±*1*.*32(*norm*)*µb*

$$
\frac{d\sigma(HF \to e)}{dy} = 92.99 \pm 3.16(stat)_{-30.08}^{+30.5}(sys)_{-11.46}^{+31.58}(extr) \pm 3.25(norm)\mu b
$$
\n
$$
\frac{d\sigma(c \to e)}{dy} = 84.52 \pm 3.50(stat)_{-32.24}^{+32.73}(sys)_{-12.15}^{+32.03}(extr) \pm 3.55(norm)\mu b
$$
\n
$$
\frac{d\sigma_c}{dy} = 889.18 \pm 36.85(stat)_{-339.16}^{+344.27}(sys)_{-127.79}^{+336.97}(extr) \pm 37.35(norm) \pm 36.74(br)\mu b
$$
\n
$$
\sigma_c = 7.59 \pm 0.31(stat)_{-2.9}^{+2.94}(sys)_{-2.49}^{+3.18}(extr) \pm 0.32(norm) \pm 0.31(br)mb
$$

dy

dy [|]BR ⁼

σ*HF*→*e*(*p*^t *>* 0*.*5*GeV/c,|y| <* 0*.*5) = 37*.*73*±*3*.*17(*stat*)

Calculation of total cross section at 2.76 TeV \sim 3 with the branching ratio of charm into electrons and the kinematical correction factor men- 552

- \bullet p_T range of measured p_T differential cross section:
- +344*.*27 +336*.*97 +3*.*9 [−]4*.*1(*sys*) +3*.*5 1< pT <8 GeV/c σ*^c* = 7*.*9*±*0*.*9(*stat*)
- = 889*.*18*±*36*.*85(*stat*) \bullet 1< p T <8 GeV/c
Extranolation based on FONNI_shane • Extrapolation based on FONNL shape
	- \bullet down to 0 p $_{\text{T}}$ for d σ (b \rightarrow e)/dy (64.3% unmeasured based \circ *dy* ∆*y* • down to 0 p_T for d σ (b \rightarrow e)/dy (64.3% unmeasured based on FONLL)
	- \bullet to full rapidity for total bp cross section or total b = 87*.*8*±*10*.*5(*stat*) $\frac{4}{3}$ • to full rapidity for total bb cross section

$$
\frac{d\sigma(b,c \to e)}{dy} = \frac{\sigma_{vis} * \beta_{b,c}}{\Delta y} \qquad \text{where} \qquad \beta_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 \text{GeV}/c < p_t < \infty, y_{min} < y < y_{max})}{\sigma_{vis, \text{FONLL}}} \\ \frac{d\sigma_{b,c}}{dy} = \frac{\sigma_{vis} * \gamma_{b,c}}{\Delta y} \qquad \text{where} \qquad \gamma_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c, 0 \text{GeV}/c < p_t < \infty, y_{min} < y < y_{max})}{\sigma_{vis, \text{FONLL}}} \\ \sigma_{b,c} = \frac{\alpha_{b,c} * \sigma_{vis}}{BR(b,c \to e)} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \text{where} \qquad \alpha_{b,c} = \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{vis, \text{FONLL}}} \qquad \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{\text{FONLL}}} \qquad \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, y_{min} < y < y_{max})}{\sigma_{\text{FONLL}}} \qquad \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{\text{FONLL}}} \qquad \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{\text{FONLL}}} \qquad \frac{\sigma_{\text{FONLL}}(b,c \to e, 0 < p_t < \infty, -\infty < y < \infty)}{\sigma_{\text{
$$

⁵⁵⁸ and for the corresponding cross section of electrons from charm hadron decays

= 42*.*25*±*3*.*47(*stat*)

dy

[−]11*.*88(*sys*)

• Statistical, Systematic and normalization uncertainties:
Note: Only CTEO6.6 is available for $\overline{}$ $\overline{}$ $\overline{}$ $\overline{}$

<mark>563 Summary 200</mark>

• scaled by multiplication factors

ultiplication factors
 1988
 1989
 1989 uncertainties: entry the total values of the integration of the integration regligible effect at 7 TeV analysis) = 87*.*8*±*10*.*5(*stat*) Note: Only CTEQ6.6 is available for

[−]2*.*23(*extr*)*±*1*.*49(*norm*)*µ^b*

- Extrapolation uncertainties: = 8*.*40*±*0*.*69(*stat*) +2*.*43 [−]2*.*36(*sys*) +0*.*19
	- · each parameter is recalculated for different variation of mass, scale in FONLL [−]0*.*40(*extr*)*±*0*.*29(*norm*)*µ^b* $\overline{}$ 5 Tecalculated for different variation of mass, seale in FONLL $\overline{}$ **or different variation of mass, scale** i
- each parameter is recalculated for unferent variation or mass, scale in Forvice
• cross section obtained after variation is compared to the cross section obtained using the central value from FONLL and quadratically summed ⁵⁵⁰ gible compared to the charm contribution for 0*.*7*GeV/c < p*^t *<* 1*GeV/c*. One obtains with this method *d*σ*b,^c*

$\overline{0}$ Total cross sections Total cross sections σ*vis,FONLL* **355 from heaviliers** and be determined at mid-rapidity to be determined at α m_{Ω} ³⁶² section obtained using the central values from FONLL. The integration range of the visible cross section

^γ*b,^c* ⁼ ^σ*FONLL*(*b,c,*0 GeV*/^c < ^p*^t *<* [∞] ³⁶⁵ Here, the range in rapidity is *|y| <* 0*.*8. One obtains for the cross section of electrons from beauty hadron <u>VISIL</u> Visible cross section

 $\mathbf{0}_{b\to e}(p_{\rm T} > 1$ Uev $\vert t \rangle \setminus \vert y \vert \leq 0.6$) $-$ 5.44 \pm 0.4. $\frac{\sqrt{3556 \text{ s}}}{\sigma_v}$ ($n_r > 1$ GeV /c |v| < 0.8) – 3.44 + 0.41(stat)^{+1.09}(svs) + 0.066(norm) ub $\sigma_{b\rightarrow e}(p_{\rm T} > 1 GeV/c, |y| < 0.8) = 3.44 \pm 0.41(stat)_{-1.30}^{+1.09}(sys) \pm 0.066(norm)\mu b$

,ymin < y < ymax)

Efficiency (2.76 TeV) ³¹⁴ Corrections are applied to the background-subtracted electron spectra for the geometrical acceptance of

The product of the overall acceptance and efficiency (εgeo [×] ^εreco [×] ^ε eID [×] ^εIP ³¹⁷) as a function of *^p*^T is